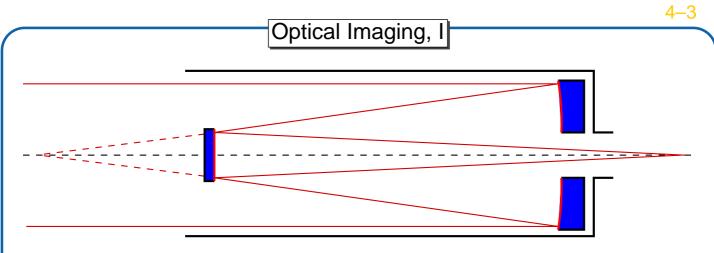


(Introduction
	Detectors we have dealt with so far: non-imaging detectors.
	This lecture: imaging optical photons and X-rays
	Any imaging system has two parts:
	 Imaging optics In most applications, mirrors are used for imaging, although other techniques (variants of shadow cameras) are used, e.g., for γ-rays
	 Spatially resolved detector a detector capable of measuring where a photon hits it in the focal plane. This can be a photographic plate or film, but is normally a charge coupled device (CCD).

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Cassegrain telescope, after Wikipedia

Reminder: Optical telescopes are usually reflectors:

primary mirror \rightarrow secondary mirror \rightarrow detector

Main characteristics of a telescope:

- collecting area (i.e., open area of telescope, $\sim \pi d^2/4$, where d: telescope diameter)
- fpr small telescopes: angular resolution,

$$\theta = 1.22 \frac{\lambda}{d} \tag{4.1}$$

but do not forget the seeing!

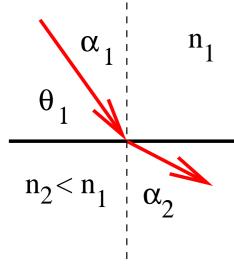
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Optical Imaging, II

Optical telescopes are based on principle that reflection "just works" with metallic surfaces. For X-rays, things are more complicated...



Snell's law of refraction:

$$\frac{\sin \alpha_1}{\sin \alpha_2} = \frac{n_2}{n_1} = n \tag{4.2}$$

where n index of refraction, and $\alpha_{1,2}$ angle wrt. surface normal. If $n \gg 1$: Total internal reflection

Total reflection occurs for $\alpha_2 = 90^\circ$, i.e. for

 $\sin \alpha_{1,c} = n \quad \Longleftrightarrow \quad \cos \theta_c = n \tag{4.3}$

with the critical angle $\theta_{c} = \pi/2 - \alpha_{1,c}$. Clearly, total reflection is only possible for n < 1

Light in glass at glass/air interface: $n = 1/1.6 \Longrightarrow \theta_c \sim 50^\circ \Longrightarrow$ principle behind optical fibers.

4–5

Optical Imaging, III

X-rays: theory gives index of refraction vacuum versus material as

$$n = 1 - N_{\mathsf{A}} \frac{Z}{A} \frac{r_{\mathsf{e}}}{2\pi} \frac{\rho}{\lambda^2} =: 1 - \delta$$
(4.4)

 $N_{\rm A}$: Avogadro's number, $r_{\rm e} = 2.8 \times 10^{-15}$ m, Z: atomic number, A: atomic weight ($Z/A \sim 0.5$), ρ : density, λ : wavelength (X-rays: $\lambda \sim 0.1-1$ nm).

Critical angle for X-ray reflection:

 $\cos\theta_{\rm c} = 1 - \delta \tag{4.5}$

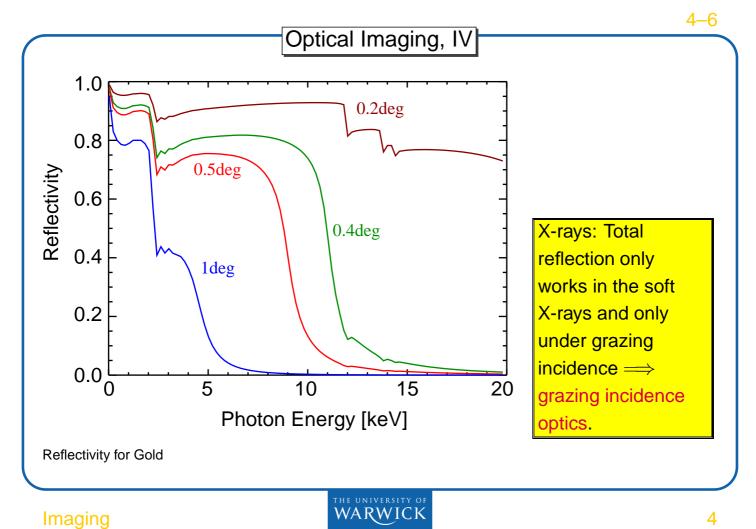
Since $\delta \ll 1$, Taylor ($\cos x \sim 1 - x^2/2$):

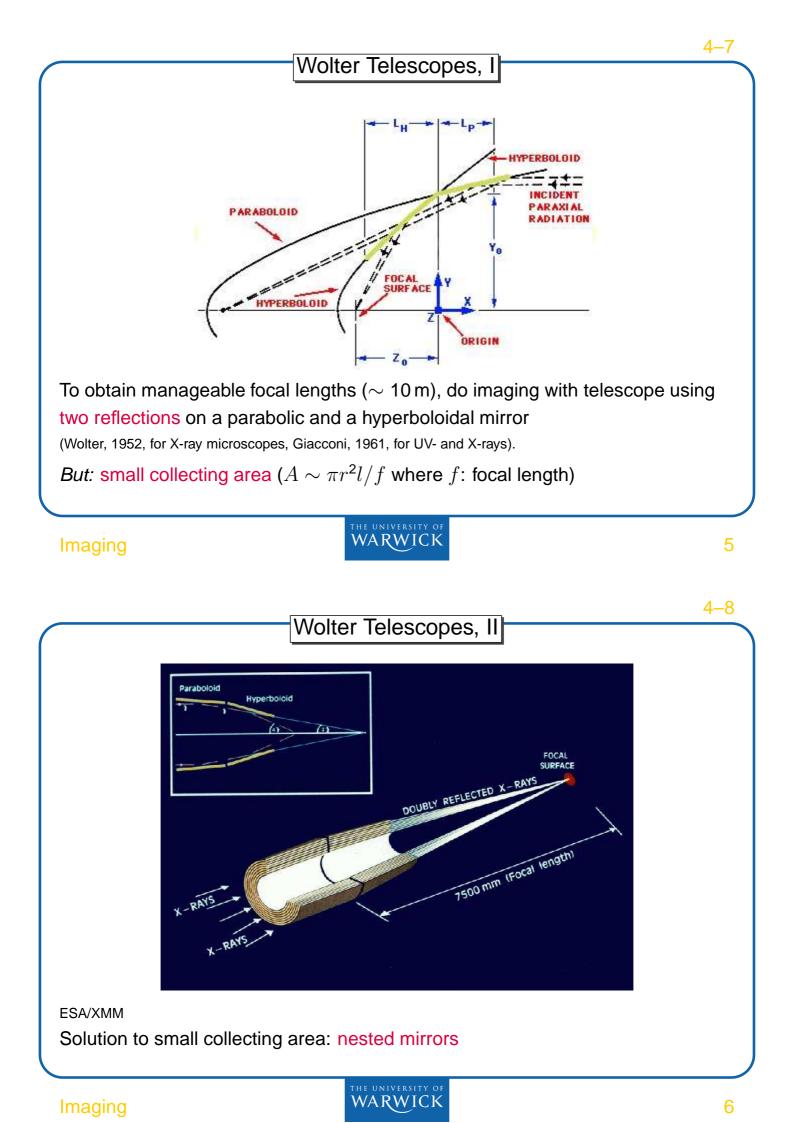
$$\theta_{\rm c} = \sqrt{2\delta} = 56' \rho^{1/2} \frac{\lambda}{1\,\rm nm} \tag{4.6}$$

So for $\lambda \sim 1$ nm: $\theta_{\rm c} \sim 1^{\circ}$.

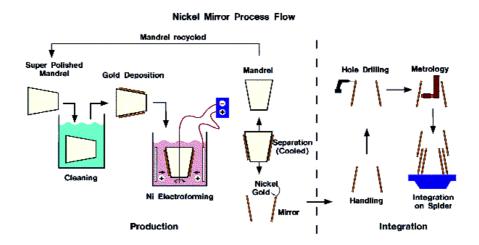
To increase θ_c : need material with high ρ \implies gold (*XMM-Newton*) or iridium (*Chandra*).

Imaging





Mirror manufacture, I



Recipe for making an X-ray mirror:

- 1. Produce mirror negative ("Mandrels"): Al coated with Kanigen nickel (Ni+10% phosphorus), super-polished [0.4 nm roughness]).
- 2. Deposit 250 nm Au onto Mandrel
- 3. Deposit 1 mm Ni onto mandrel ("electro-forming", 10 μ m/h)
- 4. Cool Mandrel with liquid N. Au sticks to Nickel
- 5. Verify mirror on optical bench.

Total production time of one mirror: 12 d, for XMM: 3×58 mirrors.

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Gold plastered mandrel for one of the XMM mirrors before electroforming the Ni shell onto the gold.

ESA picture 96.05.006-070

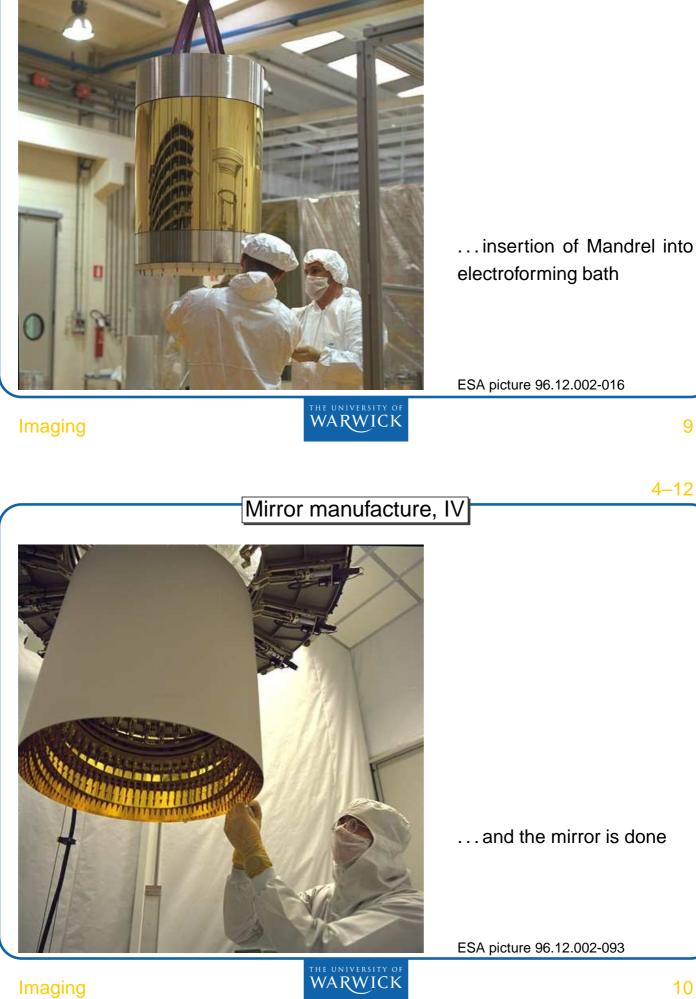
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Mirror manufacture, III



XMM-Newton



Top of the XMM mirrors: 3 mirror sets, each consisting of 58 mirrors,

- Thickness between 0.47 and 1.07 mm
- Diameter between 306 and 700 mm,
- Masses between 2.35 and 12.30 kg,
- Mirror-Height 600 mm
- Reflecting material: 250 nm Au.

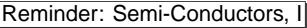
photo: Kayser-Threde

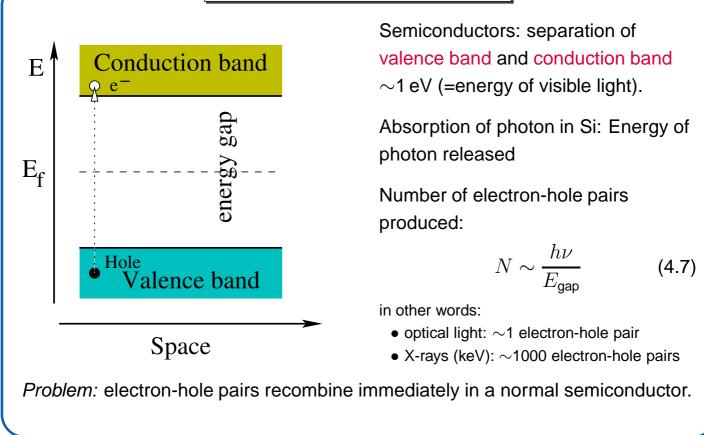
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The XMM-Newton Spacecraft (photo: ESA)





Charge Coupled Devices

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4 - 16Reminder: Semi-Conductors, II "Doping" the E Conduction band semiconductor moves the valence- and conduction **Conduction band** bands. Acceptors E_f Connecting a "n-type" and a "p-type" semiconductor 🔆 Hole Donors Valence band gives a pn-junction. Electron-hole pairs Valence band created at pn-junction will be separated by field n-type semiconductor p-type semiconductor gradient \implies electrons can then be collected in potential well away from the junction and read out.

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Reminder: Semi-Conductors, III

Number of electron-hole pairs produced determined by band gap + "dirt effects" ("dirt effects": e.g., energy loss going into bulk motion of the detector crystal ["phonons"])

Material	Z	Band gap	E/pair
		(eV)	(eV)
Si	14	1.12	3.61
Ge	32	0.74	2.98
CdTe	48–52	1.47	4.43
Hgl2	80–53	2.13	6.5
GaAs	31–33	1.43	5.2

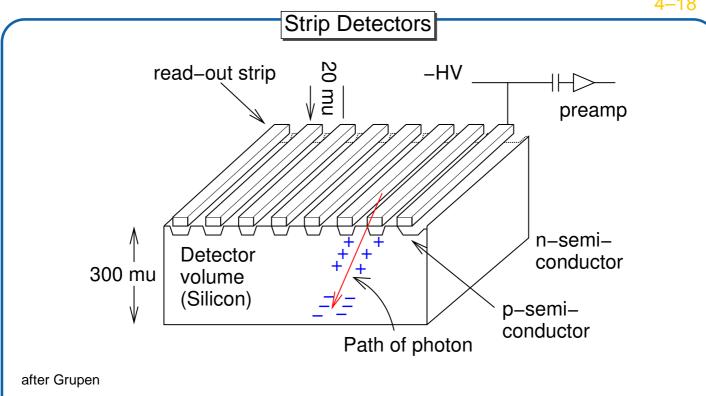
(so \sim 1 electron for optical photons, \sim 1000 electrons for X-rays)

Since band gap small: thermal noise \implies need cooling

(ground based: liquid nitrogen, -200° C, in space: cryostats)

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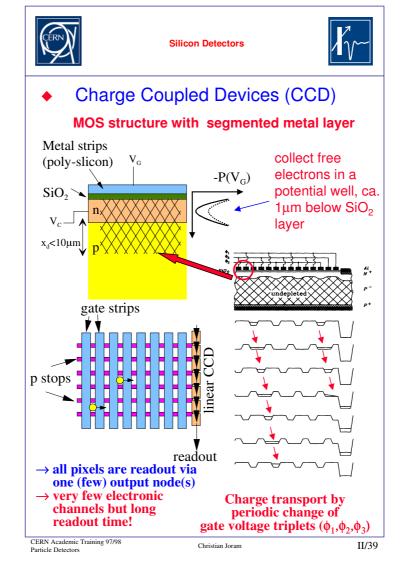


First attempt at spatial resolution obtained by segmenting the p-doted layer: microstrip detectors

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CCDs, II

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optical CCDs: mesure intensity \implies need long exposures

X-ray CCDs: measure individual photons \implies need fast readout bright sources: several 1000 photons per second \implies readout in μ s!

In X-rays: spectroscopy possible. Typical resolution reached today:

$$\frac{\Delta E}{E} = 2.355 \sqrt{\frac{F}{3.65E}} \tag{4.8}$$

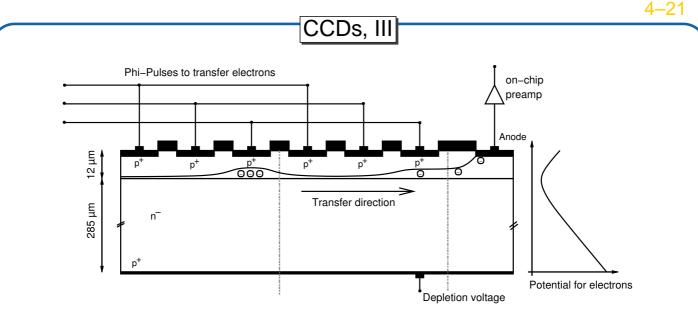
with $F \sim 0.1 \implies \sim 0.4\%$, so much better than proportional counters. (but same $\Delta E/E \propto E^{-1/2}$ proportionality because of Poisson!)

For both optical and X-rays: sensitivity close to 100%

Si based CCDs are currently the best available imaging photon detectors for optical and X-ray applications.

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Schematic structure of the XMM-Newton EPIC pn CCD.

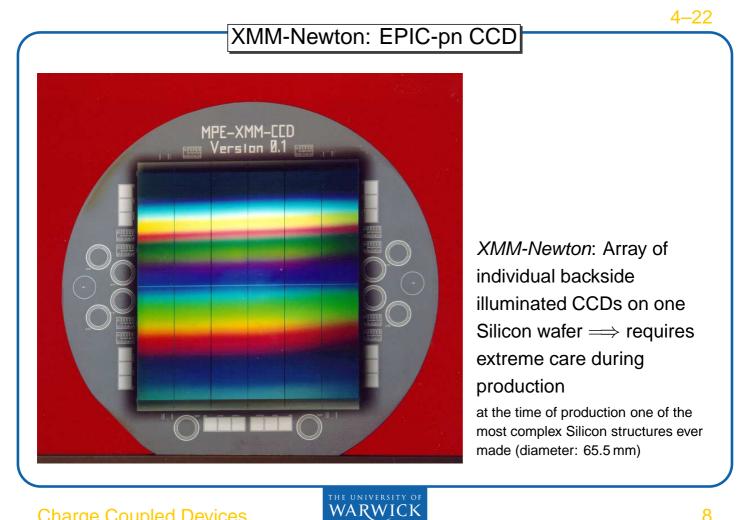
Problem: Infalling structure has to pass *through* structure on CCD surface \implies loss of low energy response, also danger through destruction of CCD structure by cosmic rays...

Solution: Irradiate the back side of the chip. Deplete whole CCD-volume, transport electrons to pixels via adequate electric field ("backside illuminated CCDs")

Note: solution works mainly for X-rays

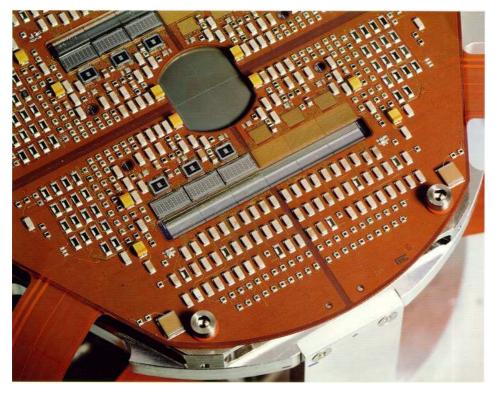
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XMM-Newton: EPIC-pn CCD



Backside of the EPIC-pn camera head

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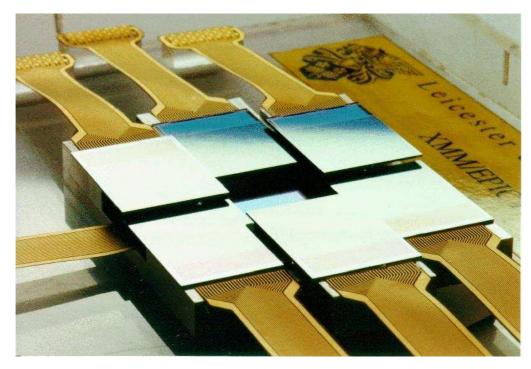
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XMM-Newton: EPIC-pn CCD



XMM-Newton (EPIC-MOS; Leicester): 7 single CCDs with 600×600 pixels, mounting is adapted to curved focal plane of the Wolter telescope.

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