



### Structure Formation and DM

Finally, the *real* linear theory has to be done in linearized or even full general relativity

⇒ very, very complicated.

**Full fledged, detailed structure formation is mainly done numerically.**

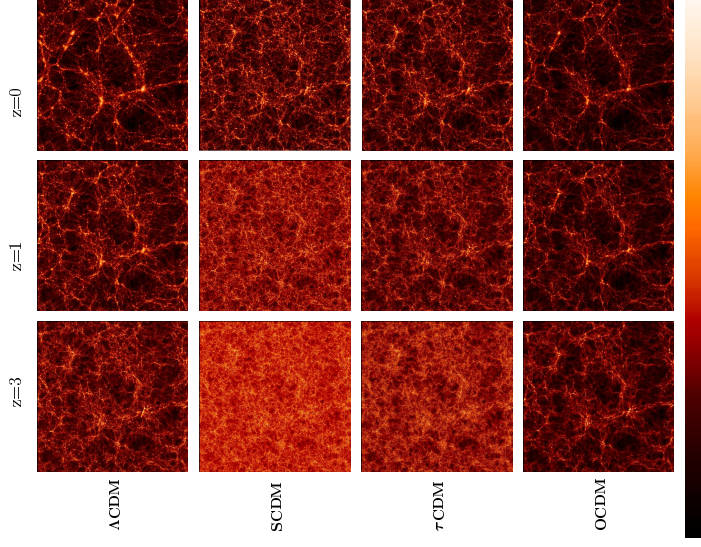
$N$ -body codes: describe particles (=galaxies) as points, compute mutual interactions in expanding universe

Requires massive computing power.

VIRGO consortium: USA, Canada, Germany, UK

Hubble Volume Simulation: Garching T3E (512 processors), 70 h CPU time followed by the Millenium Simulation (30 d CPU time)

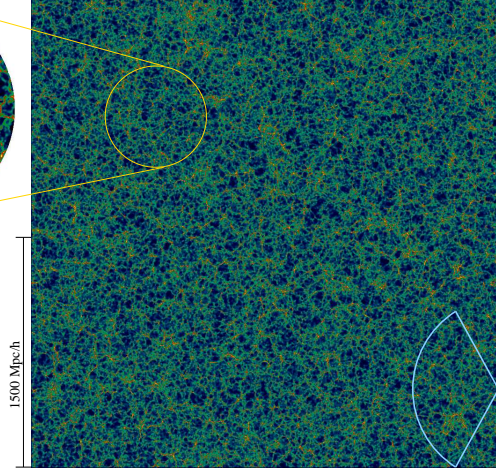
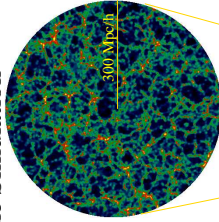
see Springel et al. (2005), Springel, Frenk & White (2006) and <http://www.mpa-garching.mpg.de/~virgo/virgo/>



### Structure Formation

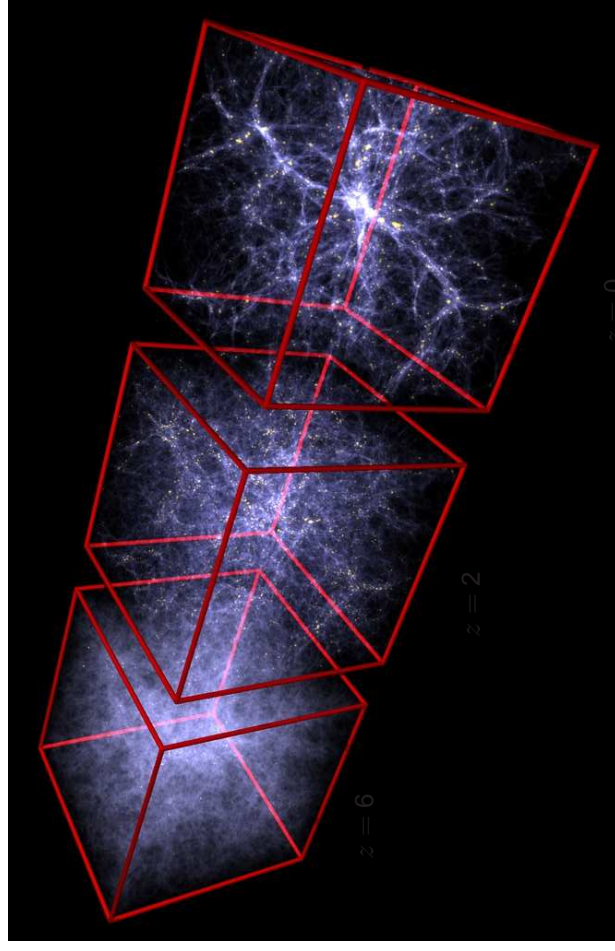
#### The Hubble Volume Simulation

$\Omega=0.3, \Lambda=0.7, h=0.7,$   
 $\sigma_8=0.9$  ( $\Lambda$ CDM)  
 $3000 \times 3000 \times 30 \text{ h}^3 \text{ Mpc}^3$   
 $P^{\text{M}}, z_i=35, \quad \bar{v}=100 \text{ h}^{-1} \text{ kpc}$   
 $1000^3$  particles,  $1024^3$  mesh  
 T3E(Garching) - 512cpus  
 $M_{\text{particle}} = 2.2 \times 10^{12} h^{-1} M_{\text{sol}}$



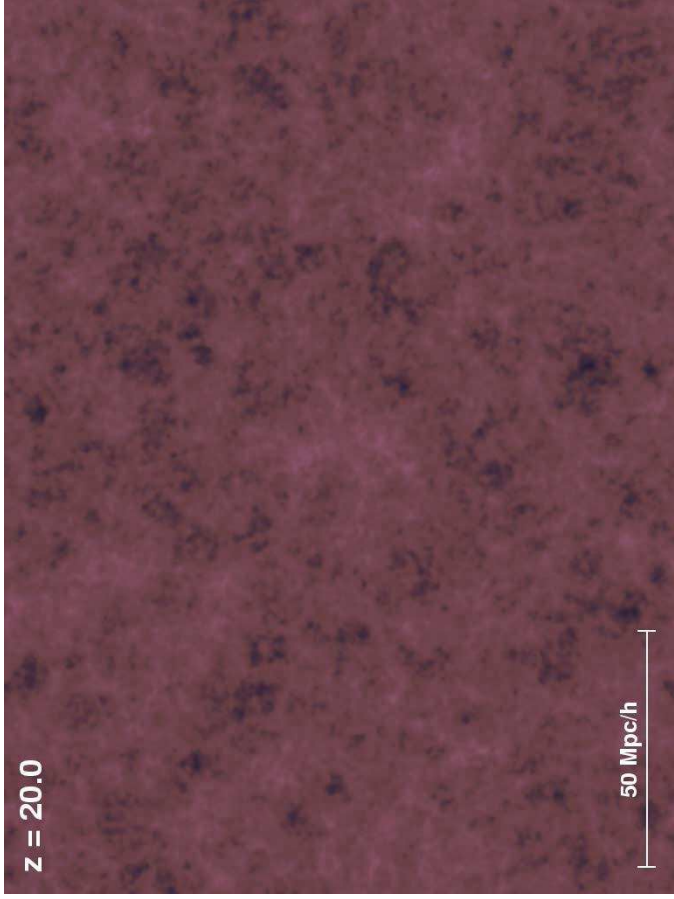
$\Lambda$ DM, pie shows SDSS size

The VIRGO Collaboration 1996

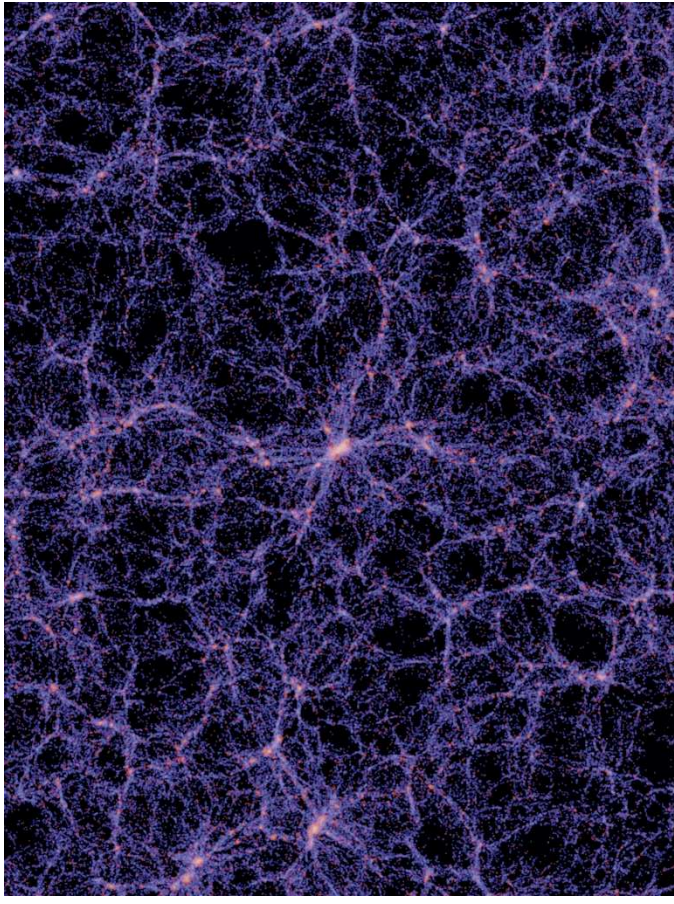


(V. Springel/MPA)

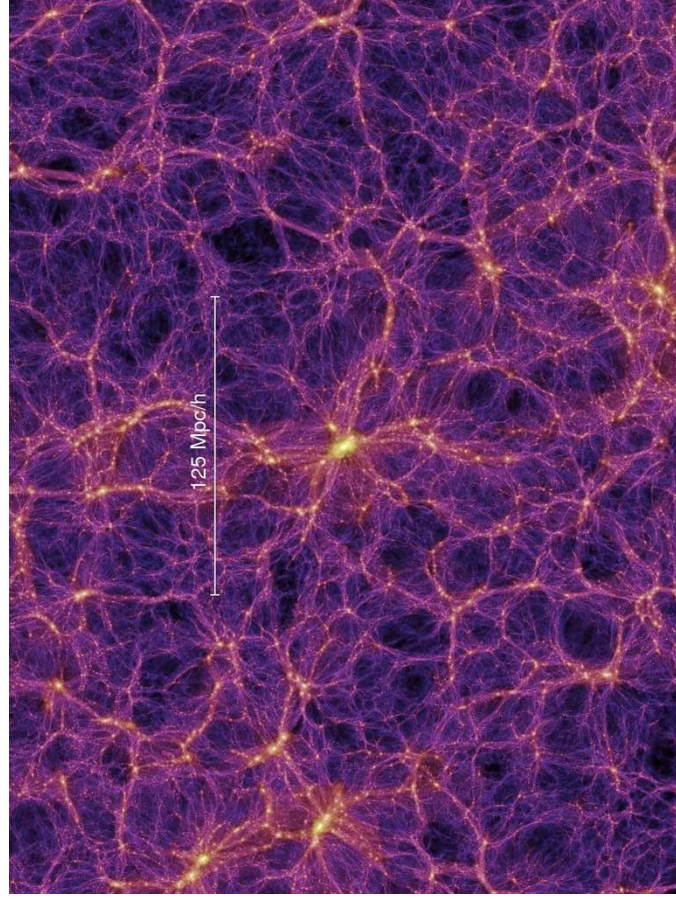




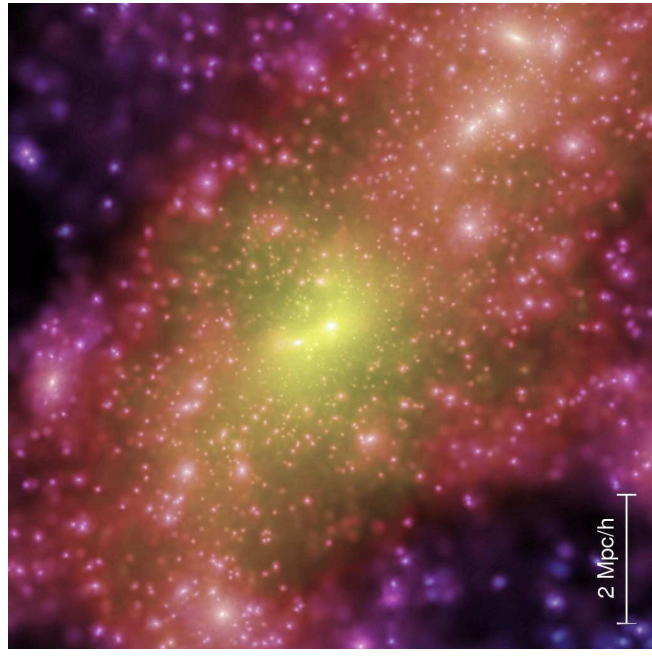
Evolution of structure in a  $\Lambda$ CDM Universe (MPA/V. Springel)



... and corresponding galaxy distribution (V. Springel/MPA)

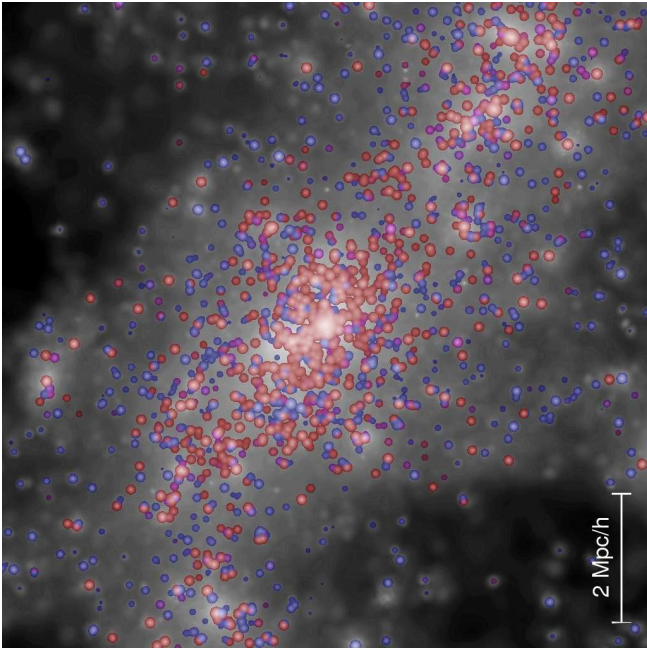


Today's dark matter distribution ... (V. Springel/MPA)

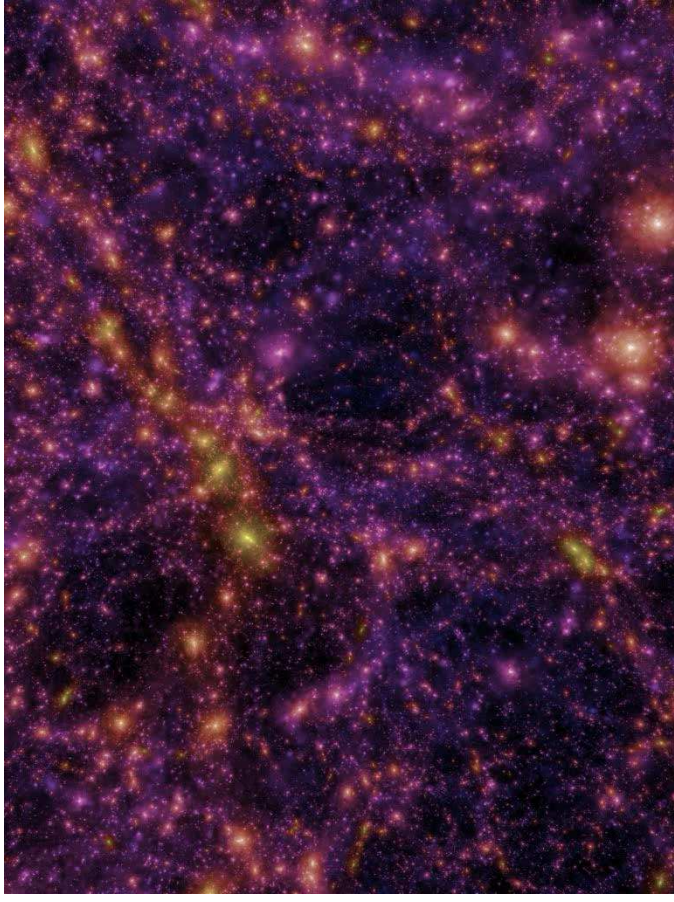


Today's dark matter distribution in a cluster ... (V. Springel/MPA)

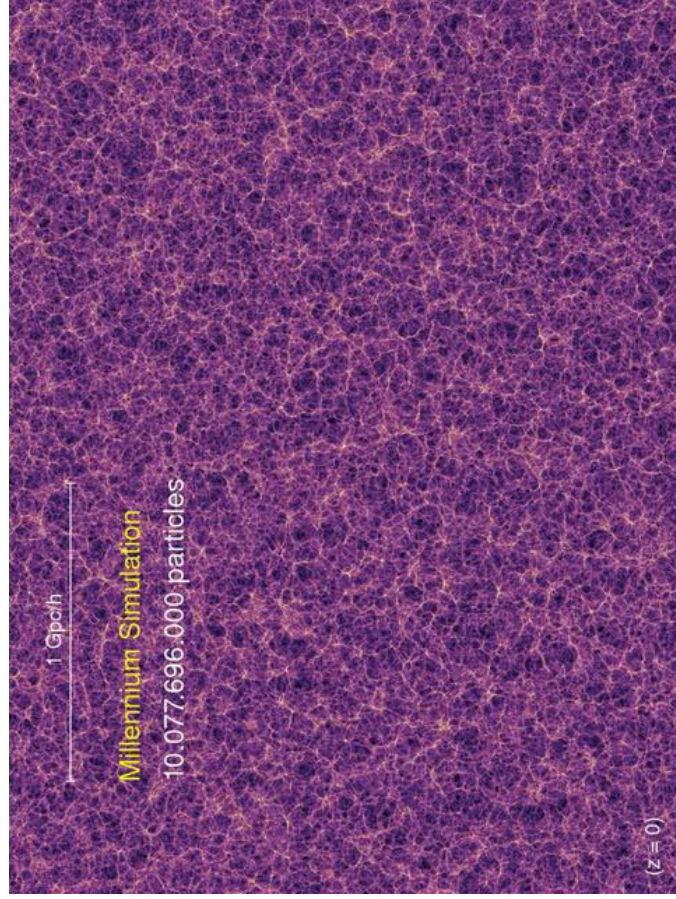




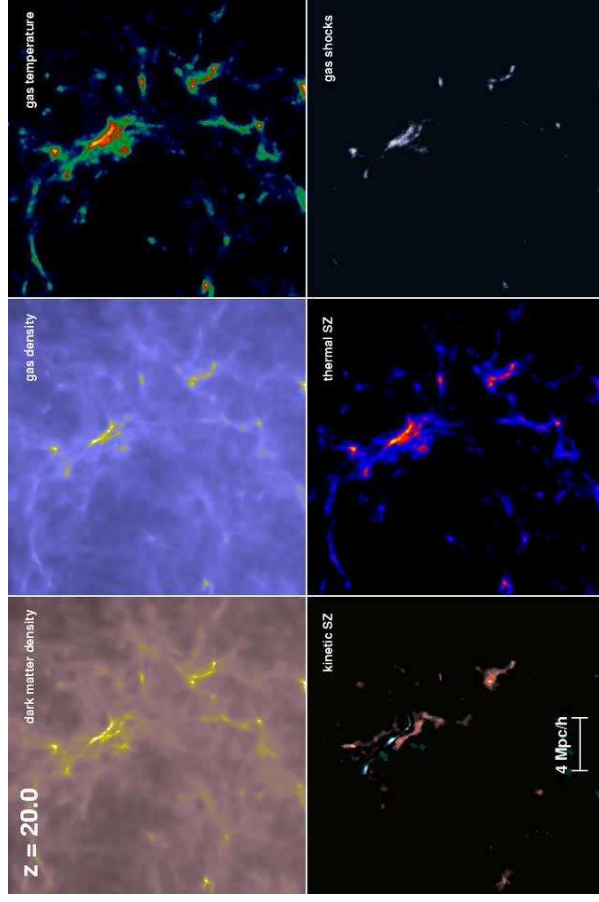
... and corresponding galaxy distribution (V. Springel/MPA)



Flight through the DM structure of the Millennium Simulation (V. Springel)



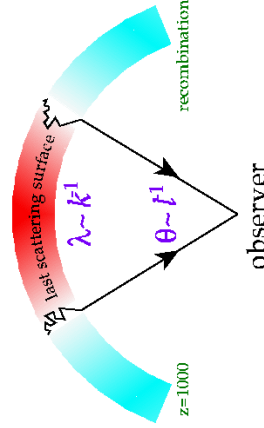
Zoom into the DM structure of the Millennium Simulation



Formation of a galaxy cluster (V. Springel)



CMBR



Matter and Radiation are coupled, i.e., large mass density = high photon density.

Photons from overdense regions: gravitational redshift  $\implies$  observable! (Sachs Wolfe Effect)

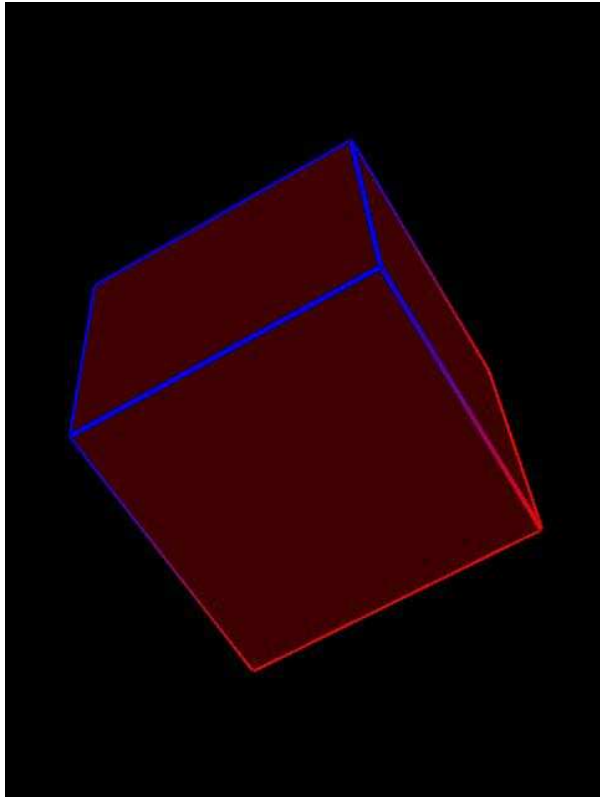
CMBR: Radiation from surface of last scattering  $\implies$  Gravitational redshift is observable as temperature fluctuation:

$$\frac{\Delta T}{T} \sim \frac{\Delta \Phi_g}{c^2} \tag{9.81}$$

CMBR Fluctuations trace gravitational potential at  $z \sim 1100!$

courtesy Wayne Hu

Initial conditions

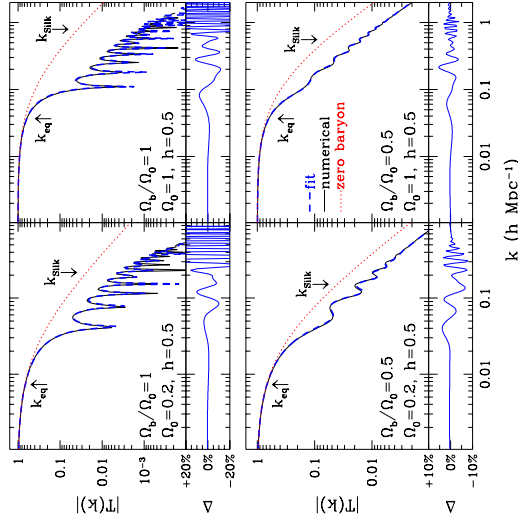


(movie courtesy N. Gnedin)

Reionization: ionization of intergalactic medium through UV radiation from the first stars



Formal Structure Formation



Calculation of real power spectra difficult: growth under self-gravitation pressure effects dissipation. To predict observations from today: define transfer function

$$\delta_k(z=0) = D(z) T_k \delta_k(z) \tag{9.80}$$

But: need initial conditions,  $\delta_k(z)!$

(Eisenstein & Hu, 1999)

Initial conditions

CMBR

Temperature fluctuations:

$$\frac{\Delta T}{T} \sim \frac{\Delta \Phi_g}{c^2} \tag{9.81}$$

where

$$\Delta \Phi_g \sim -\frac{2G\Delta M}{R} = \frac{8\pi G}{3} \bar{\rho} R^2 \delta = -\delta(t) (H(t)R)^2 \tag{9.82}$$

Current angle of region on sky:

$$\alpha \sim R/d_A \tag{9.83}$$

where the angular diameter distance

$$d_A = d_L / (1+z)^2 \tag{9.84}$$

Therefore:

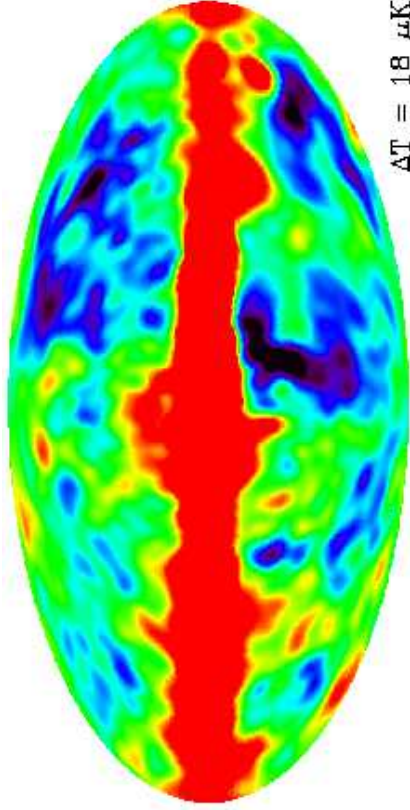
$$\frac{\Delta T}{T} \sim \frac{\Delta \Phi_g}{c^2} \propto \frac{\delta \alpha^2}{3} \tag{9.85}$$

Quotient 3 from more detailed theory, "Integrated Sachs Wolfe effect"

Initial conditions



CMBR



$$\Delta T = 18 \mu\text{K}$$

COBE: Resolution  $\alpha \sim 7^\circ$  (corresponds to  $\sim 10^{20} M_\odot$  at recombination). Temperature fluctuations imply  $\delta \sim 10^{-3}$  at recombination. This is small for pure matter dominated universe  $\implies$  Implies existence of dark matter!

Initial conditions

CMBR

Detailed theory of fluctuations: Expand CMB fluctuations on sky in spherical harmonics:

$$\frac{\Delta T}{T}(\theta, \phi) = \sum_{\ell, m} a_{\ell, m} Y_{\ell, m}(\theta, \phi) \tag{9.86}$$

Since rotationally symmetric, can express variation in terms of multipole coefficients,  $C_\ell$ :

$$C(\theta) = \frac{1}{4\pi} \sum_{\ell} \sum_{m=-\ell}^{+\ell} |a_{\ell, m}|^2 P_\ell(\cos\theta) \tag{9.87}$$

$$=: \frac{1}{4\pi} \sum_{\ell} (2\ell + 1) C_\ell P_\ell(\cos\theta) \tag{9.88}$$

where  $C(\theta) = \langle \Delta T/T \rangle$  and where the  $P_\ell$  are the Legendre polynomials.

Initial conditions

CMBR

Expect following features:

**Large angle anisotropy:** (small  $\ell$ , scales  $\gtrsim$  horizon at decoupling): Flat part due to Sachs-Wolfe effect

**Smaller angular scales:** (larger  $\ell$ ): Influenced by photon-baryon interactions:

Matter falls in potential well

$\implies$  Pressure resists

$\implies$  acoustic oscillations

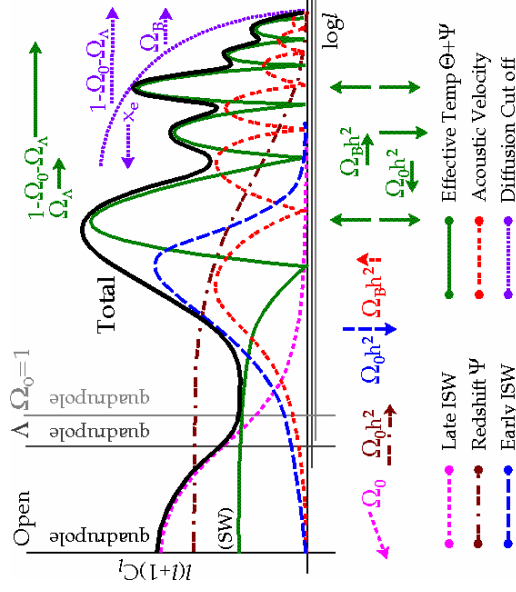
$\implies$  Power at selected scales!

Power from those density fluctuations which had their maximum amplitude at time of last scattering dominates  $\implies$  acoustic peak

Also damping from photon diffusion (Compton scattering; Silk damping [after Joseph Silk])

Initial conditions

CMBR

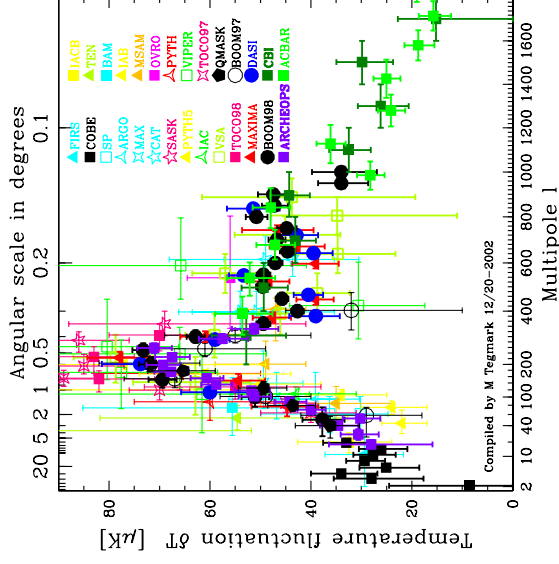


Hu, Sugiyama, & Silk (1995)

Initial conditions



Summary: Pre-WMAP



Power Spectrum of CMB

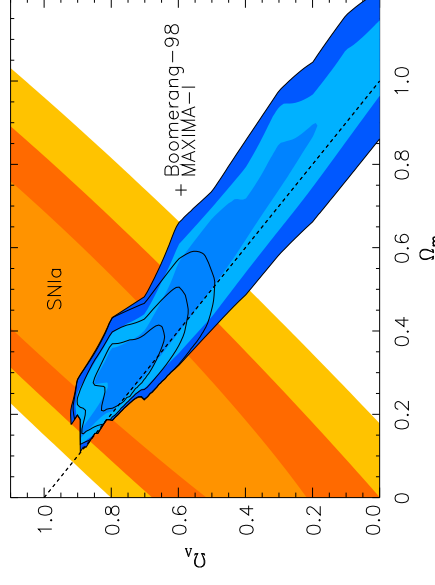


Enter: BOOMERANG (Balloon Observations of Millimetric Extragalactic Radiation and Geophysics), Flight in Antarctica 1998 December 29 – 1999 January 9



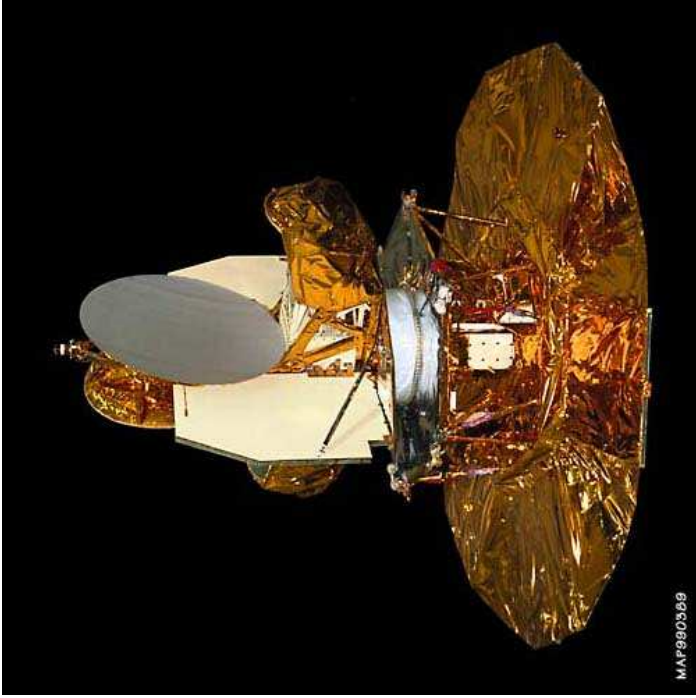
BOOMERANG before Mt. Erebus; courtesy BOOMERANG team  
Other balloon missions: MAXIMA-1, ...

Summary: Pre-WMAP

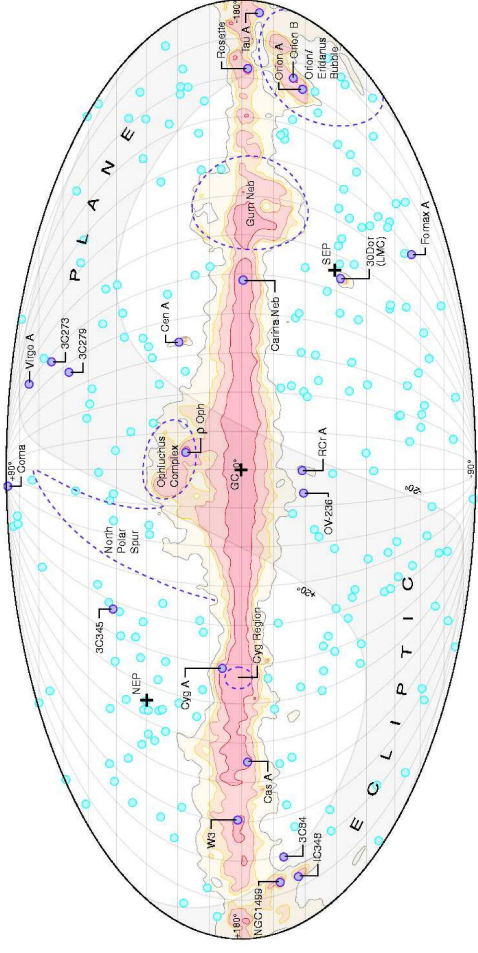
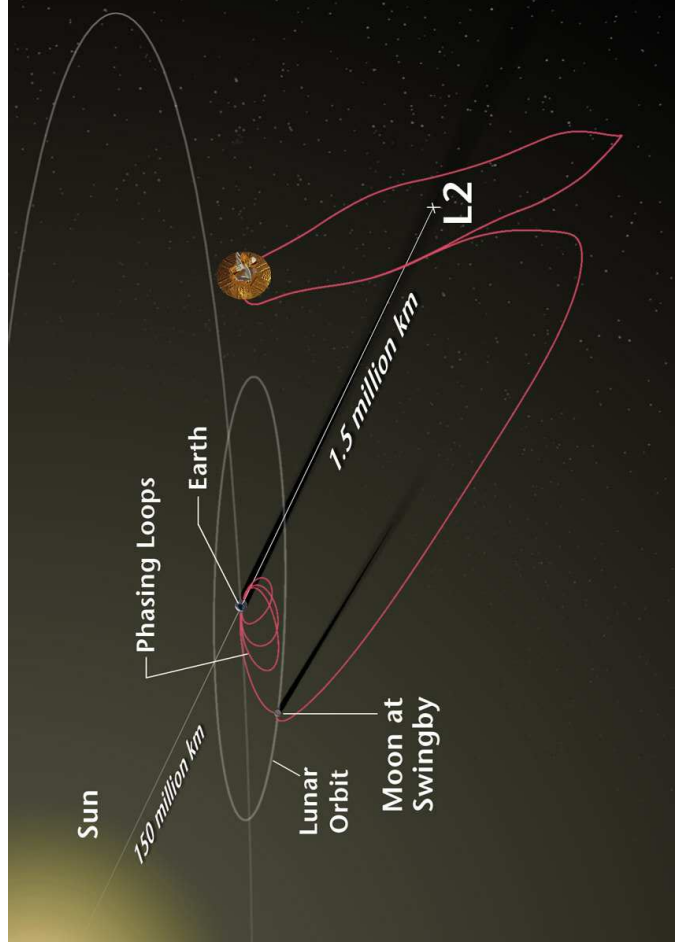


Power Spectrum of CMB



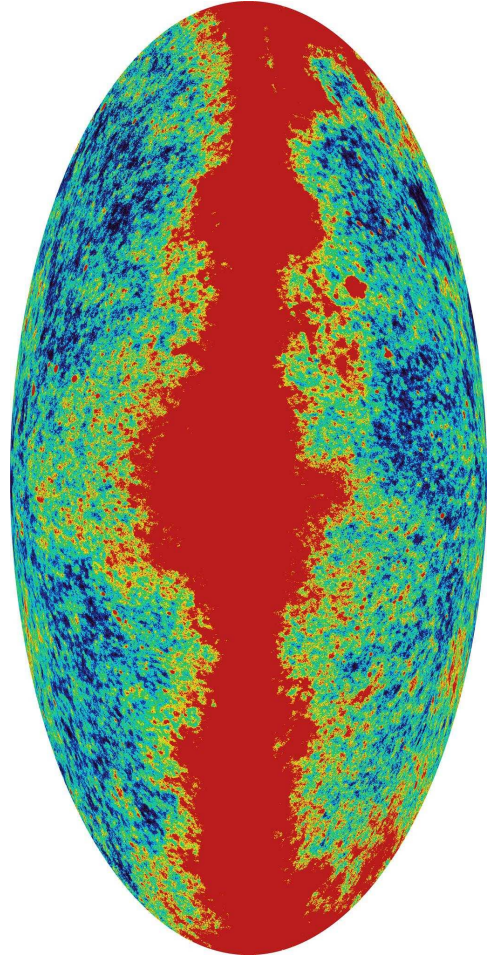


- Wilkinson Microwave Anisotropy Probe (WMAP):
- Launched 2001 June 30, measurements began 2001 August 10
  - Orbit around 2nd Lagrange Point of Sun-Earth System
  - Highly precise radiometers of high spatial resolution (best:  $0.21^\circ$  FWHM in W-Band at 3.2 mm) in five wavebands (see Bennett et al. 2003 for an overview).



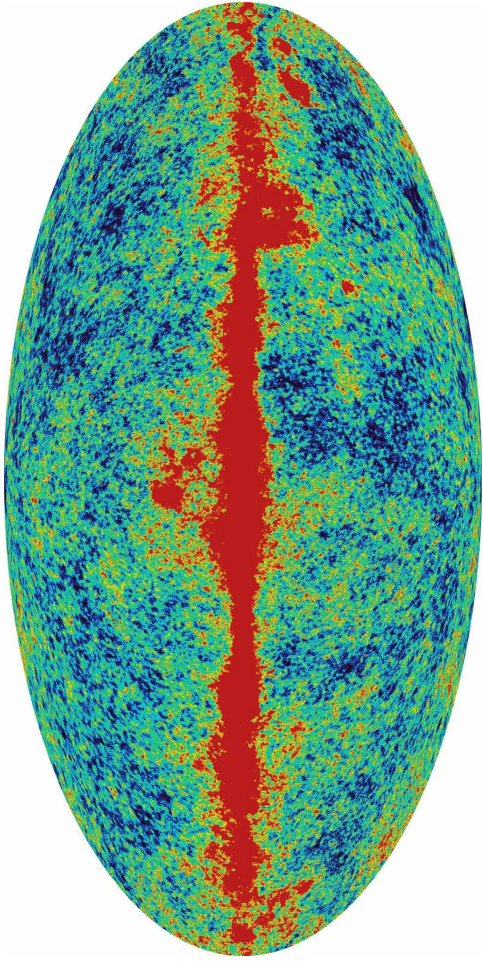
**Foreground features of the microwave sky (Bennett et al., 2003).**

Sunyaev Zeldovich effect is expected to be strongest in Coma cluster, temperatures of  $-0.34 \pm 0.18$  mK in W and  $-0.24 \pm 0.18$  mK in K-band; barely detectable with WMAP, does not contaminate maps.

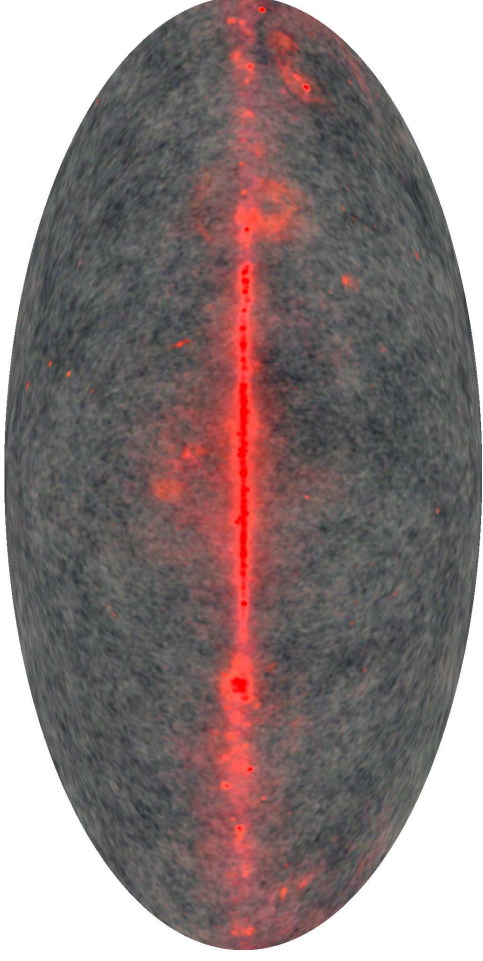


WMAP, K-Band,  $\lambda = 13$  mm,  $\nu = 22.8$  GHz,  $\theta = 0.83^\circ$  FWHM





WMAP, Q-Band,  $\lambda = 7.3 \text{ mm}$ ,  $\nu = 40.7 \text{ GHz}$ ,  $\theta = 0.49^\circ$  FWHM

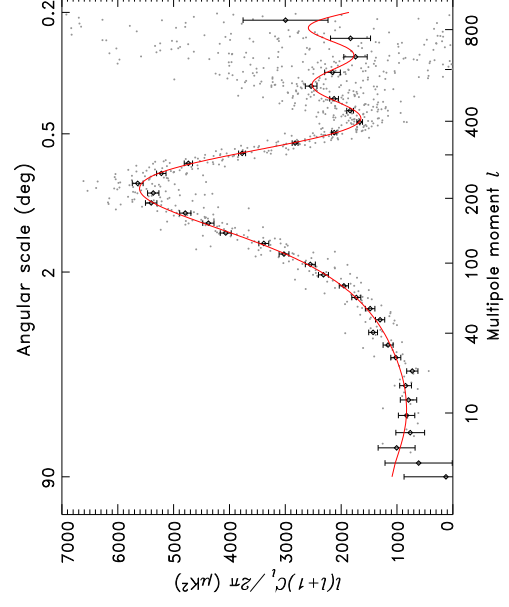


Different spectral signature enables identification of Galaxy foreground radiation



9-91

### Power Spectrum, $l$



Best fit parameters for WMAP data after 3 years of measurements and assuming  $\Omega = 1$  (Spergel et al., 2007):

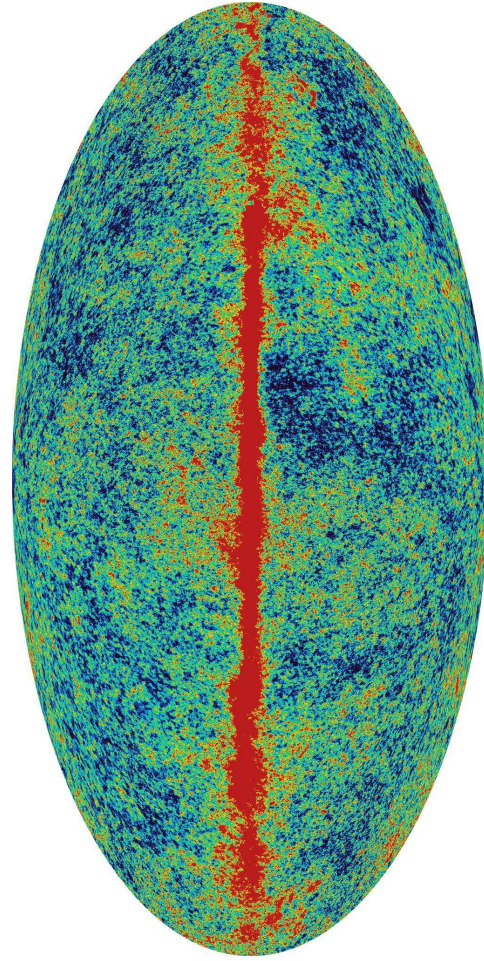
$$\begin{aligned}
 h &= 0.732 \pm 0.032 \\
 \Omega_m h^2 &= 0.1277 \pm 0.0080 \\
 \Omega_b h^2 &= 0.02229 \pm 0.00073 \\
 \tau &= 0.089 \pm 0.030
 \end{aligned}$$

and initial density fluctuations with an amplitude of  $\sigma_8 = 0.761 \pm 0.049$  and slope  $0.958 \pm 0.016$ .

Power spectrum requires that  $\Lambda$  behaves like a cosmological constant.

⇒ Very good agreement between data and theory

(WMAP, 1 year data Spergel et al., 2003, Fig. 1)



WMAP, W-Band,  $\lambda = 3.2 \text{ mm}$ ,  $\nu = 93.5 \text{ GHz}$ ,  $\theta = 0.21^\circ$  FWHM





## Power Spectrum, II

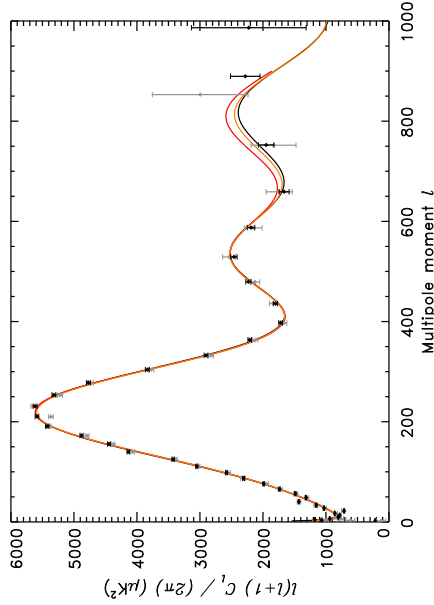
Best fit parameters for WMAP data after 3 years of measurements and assuming  $\Omega_c = 1$  (Spergel et al., 2007):

$$\begin{aligned} h &= 0.732 \pm 0.032 \\ \Omega_m h^2 &= 0.1277 \pm 0.0080 \\ \Omega_b h^2 &= 0.02229 \pm 0.00073 \\ \tau &= 0.089 \pm 0.030 \end{aligned}$$

and initial density fluctuations with an amplitude of  $\sigma_8 = 0.761 \pm 0.049$  and slope  $0.958 \pm 0.016$ .

Power spectrum requires that  $\Lambda$  behaves like a cosmological constant.

$\Rightarrow$  Very good agreement between data and theory



(WMAP, 3 year data Spergel et al., 2007, Fig. 2)

Bahcall N. A., & Soneira, R. M., 1983, ApJ, 270, 20

Bennett C. L., et al., 2003, ApJ, submitted

Cole, S., et al., 2005, MNRAS, 362, 505

de Lapparent, V., Geller, M. J., & Huchra, J. P., 1986, ApJ, 302, L1

Eisenstein, D. J., & Hu, W., 1999, ApJ, 511, 5

Hamilton, A. J. S., & Tegmark, M., 2002, MNRAS, 330, 506

Jaffe, A. H., et al., 2000, Phys. Rev. Lett., submitted (astro-ph/0007333)

Peacock, J. A., 1999, Cosmological Physics, (Cambridge: Cambridge Univ. Press)

Peebles, P. J. E., 1980, The Large-Scale Structure of the Universe, (Princeton, NJ: Princeton Univ. Press)

Spergel, D. N., et al., 2007, Astrophys. J., Suppl. Ser., 170, 377

Spergel, D. N., et al., 2003, ApJ, submitted

Springel, V., Frenk, C. S., & White, S. D. M., 2006, Nature, 440, 1137

Springel, V., et al., 2005, Nature, 435, 629

Strauss, M. A., 1999, in Structure Formation in the Universe, ed. A. Dekel, J. P. Ostriker, (Cambridge: Cambridge Univ. Press)

Strauss, M. A., & Willick, J. A., 1995, Phys. Rep., 261, 271

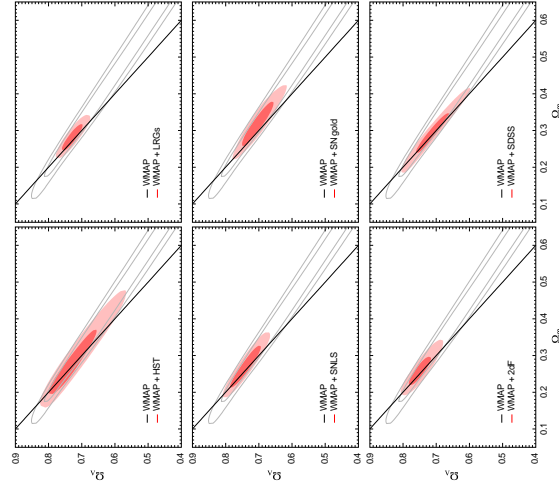
Tegmark, M., et al., 2004, ApJ, 606, 702

Tucker, D. L., et al., 1997, MNRAS, 285, L5

WMAP



## Omega



Using WMAP data alone gives

$$\Omega_m = 0.415 \quad \Omega_\Lambda = 0.630 \quad (9.90)$$

Combining WMAP with other

measurements generally results in

$\Omega_k \sim -0.01 \pm 0.01$ , i.e., measurements consistent with a flat universe.

In the year 1 data, a model with  $\Omega_\Lambda = 0$  is found to be consistent with the WMAP data only if  $H_0 = 32.5 \text{ km s}^{-1} \text{ Mpc}^{-1}$  and  $\Omega_{\text{tot}} = 1.28$   
 $\Rightarrow$  Ruled out by other measurements.

(Spergel et al., 2007, Fig. 21)

WMAP

# The End