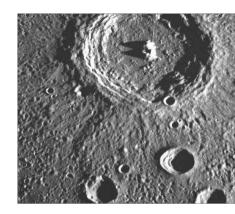
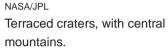
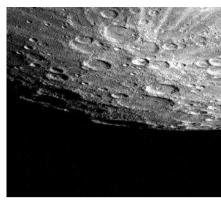
6-1

Planets: Surfaces and Interiors

Major landforms: Craters







S-Pole; NASA/JPL 50 km diam craters with rays (remains from impact)

Mercury:

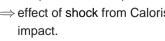


Caloris Basin (1300 km diameter) close to sub-solar point at perihelion $\Longrightarrow \text{hot } (T>400^{\circ}\text{C on day,} \\ T\sim -170^{\circ}\text{C during night)}$ result of large impact event



Robinson, NWU / NASA
Hilly/lineated terrain antipodal to
Caloris (120 km across)

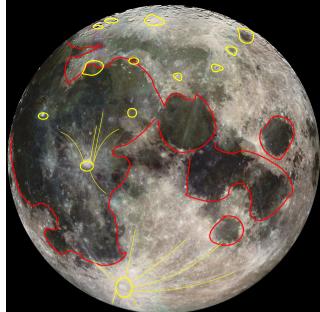
⇒ effect of shock from Caloris



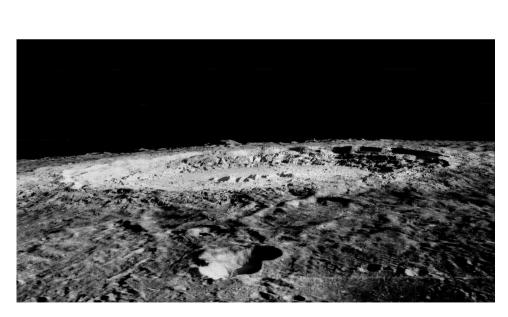


V. L. Sharpton

Earth: Wolf Creek Crater, Australia Currently 172 confirmed impact structures on Earth



Earth's Moon: surface dominated by mariae (large, dark lava basins) and craters (only most prominent shown).



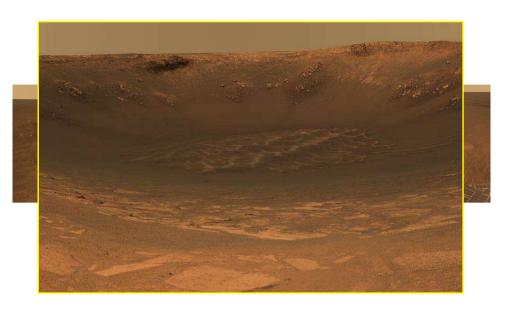
Moon: Crater Copernicus



Moon: Apollo 16, 1972 Apr, Descartes Highlands



Mars: Surface panorama, Exploration Rover "Opportunity" looks back to lander (2004 Feb 09)

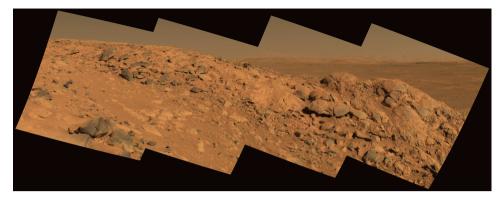


NASA/JPL/Cornell

Mars: Crater Endurance



Mars: "Spirit" rolls towards Columbia Hills (2004 June)

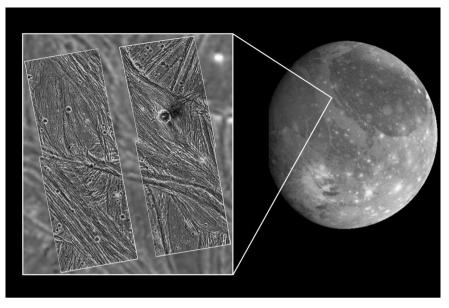


Mars: "Spirit" looks from Columbia Hills towards Gusev crater (2004 Aug)



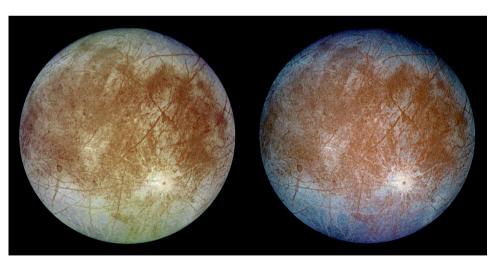
Montage of Jupiter and Galilean Moons: top to bottom: Io, Europa, Ganymede and Callisto.

(N.B.: All Galilean moons tidally locked to Jupiter – always same side is facing Jupiter)



NASA Galileo / DLR, inset: 120×110 km

Ganymede – icy surface, ice hills and valleys, craters Radius: 2634 km (~ Mercury!)



NASA Galileo / DLR, 1996 September 7

Europa – icy surface with ridges (colors: different kinds of ice)
Radius: 1565 km (~ Earth Moon)
possibility of water ocean below surface



Callisto: "pock faced", mainly impact craters.

white: ice

dark: ice-poor material

Radius: 2406 km (similar

to Mercury!)



Impact Craters

6-22

Physics of impact cratering:

Kinetic energy:

$$E = \frac{1}{2}mv^2 = \frac{1}{2} \cdot \frac{4}{3}\pi r^3 \rho v^2 = \frac{\pi d^3 \rho v^2}{12}$$

Important numbers:

- Velocity of impact: several times orbital speed of planet
- Impacting body: rock or Fe, several meters to kilometers in size

Example:

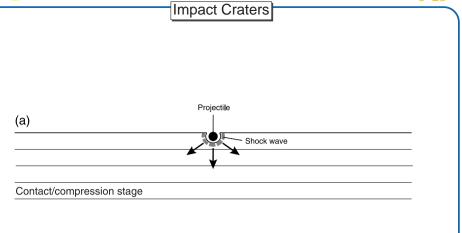
E.g.,
$$v=$$
 10 km s $^{-1}$, $d=$ 25 m, $\rho=$ 7900 kg m $^{-3}$ $\Longrightarrow E=$ 3 \times 10 15 J (\sim 1 Megaton of TNT)

1 Megaton TNT is typical strength of US nuclear bombs [B-83 bomb]

Surfaces: Craters 21



6–23



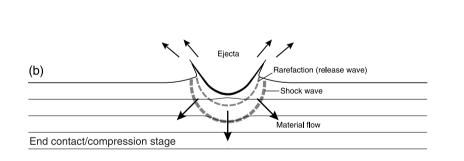
French, 1998, LPI Cont. 954

Surfaces: Craters 22

Impact Craters



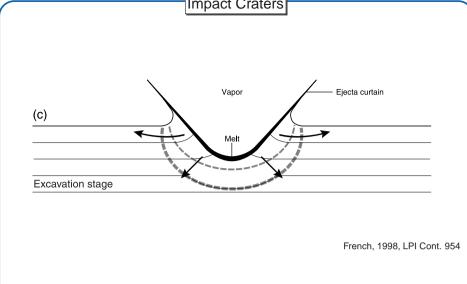
6-23



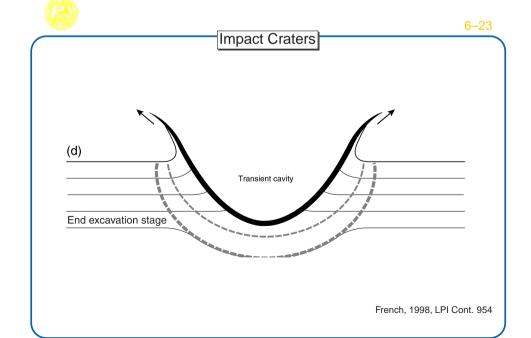
French, 1998, LPI Cont. 954

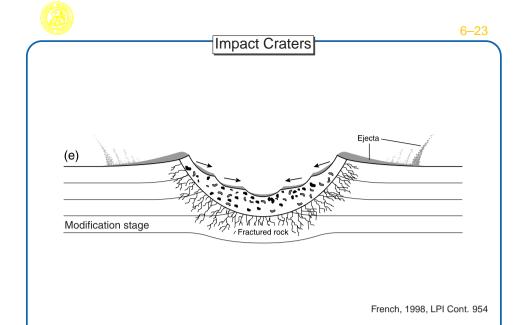
23



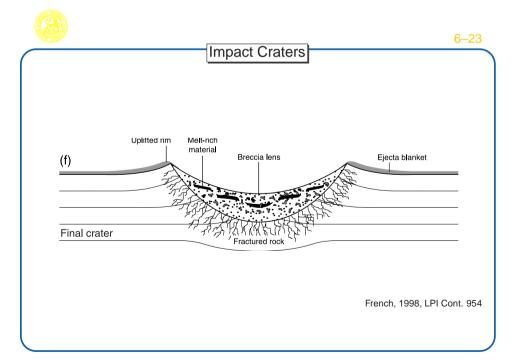


Surfaces: Craters 24

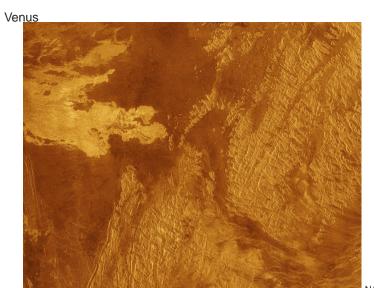




Surfaces: Craters 26



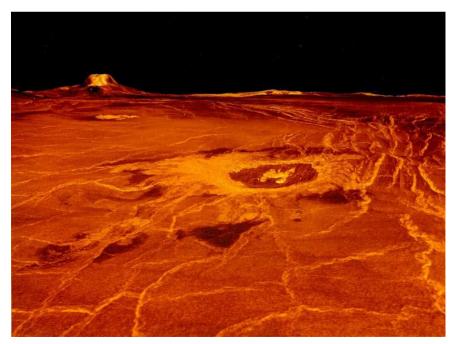
Surfaces: Craters 27



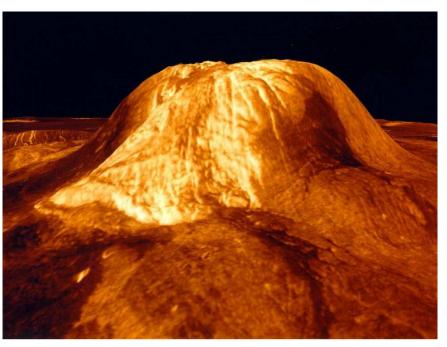
NASA, Magellan $440 \times 350 \, \text{km}^2$ area in Eistla Regio, shows basic stratigraphy (sequence of geologic events):

right half: old highlands, fractured structure (\sim 15% of surface), left part: lowlands, younger area, origin in former volcanism?

Craters (note: strong erosion ⇒ fewer craters overall)



Eistla Regio; heights exagerrated by factor 22.5

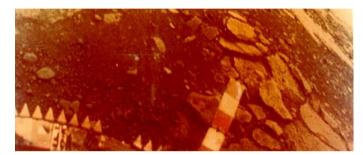


Gula Mons; heights exagerrated by factor 22.5

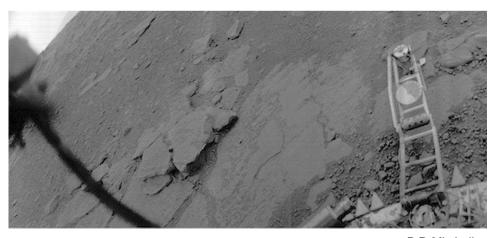
Gula Mons; real heights

Venus surface images:





Venera 13 (3 March 1982): images from color TV camera

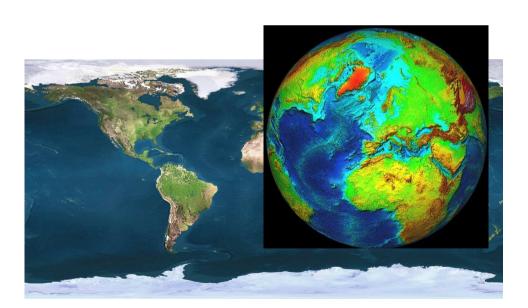


courtesy D.P. Mitchell

Venera 13 (3 March 1982): reanalysed image without camera distortion



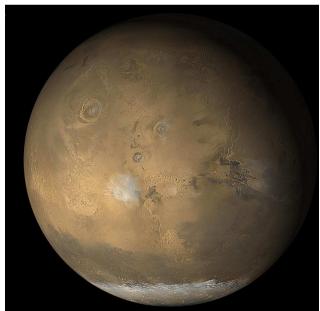
Venera 14 (5 May 1982)



Evidence for plate tectonics (few craters!)

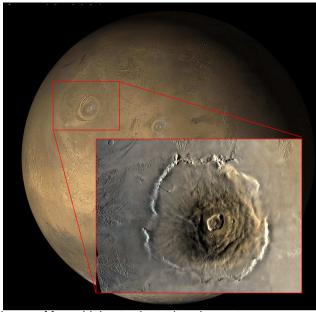


Evidence for plate tectonics (few craters!), volcanism,...

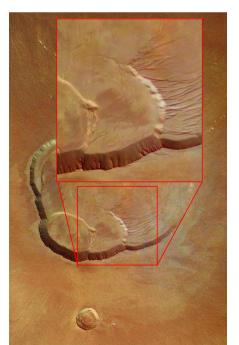


Mars: Tharsis vulcanos: Large shield vulcanos, now extinct

⇒ no plate tectonics ⇒ Mars interior is colder than Earth.



Olympus Mons: highest volcano in solar system (25 km above surrounding plain; but slope only 2° to 5°).



ESA/Mars Express, HRSC, 11.02.2004

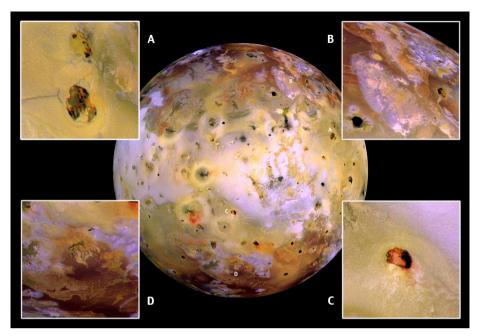


ESA/Mars Express, HRSC, 11.02.2004

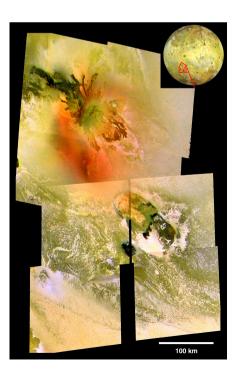


Montage of Jupiter and Galilean Moons: top to bottom: Io, Europa, Ganymede and Callisto.

(N.B.: All Galilean moons tidally locked to Jupiter – always same side is facing Jupiter)



Jupiter's moon lo – the vulcano moon (Diam. 1821 km [Earth moon: 1738 km])



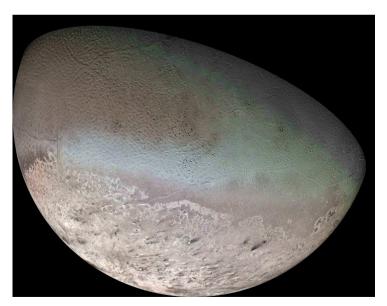
Active vulcanoes on lo (interior heated by tidal forces from Jupiter), color due to large contents of sulphur and sulphur oxides in lava.

Height of vulcanoes: 6 km or higher



curtains of lava fountains [white: overexposed]
NASA Galileo, 1999 Nov 26

High temperature volcanism (2000 K; hotter than on Earth [1700 K]!)



NASA/Voyager 2/Calvin J. Hamilton

Neptune's Moon Triton:

ice cap of frozen methane (freezing point 90 K) and frozen nitrogen (freezing point 60 K). Few impact craters \Longrightarrow young surface \Longrightarrow volcanism (dark spots: nitrogen geysers with $T\sim$ 70 K)

Interiors: Terrestrial Planets, IV

Lithosphere
Mantle
Core

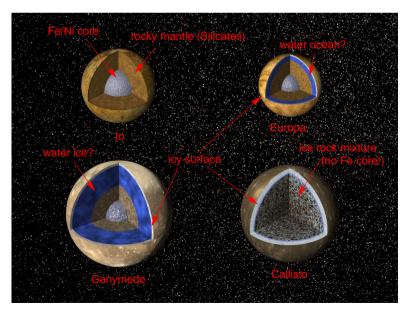
Structure of terrestrial planets:

- Core: high-density material (Fe)
- Mantle: plastic materials, hot (e.g., Earth: molten rocks)
- Lithosphere: rigid material, e.g., Silicates

Knowledge of structure important for, e.g.,

- atmospheric composition (molten mantle \Longrightarrow volcanism \Longrightarrow CO₂, CH₄,...)

Interiors



Structure of Jupiter's Galilean Moons similar to terrestrial planets (but some also have very thick ice layer on top)





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Structure: Gas Giants

In general, gas giants have very different properties from terrestrial planets:

- average density low, e.g.,
 - Jupiter: $\langle
 ho
 angle \sim$ 1.3 g cm $^{-3}$
 - Saturn: $\langle
 ho
 angle \sim$ 0.7 g cm $^{-3}$

(compare to terrestrial planets: $\langle \rho \rangle \sim 5.5\,\mathrm{g\,cm^{-3}}$; water has $\rho = 1\,\mathrm{g\,cm^{-3}}$).

- elemental composition similar to stars (by mass):
 - 75% H
 - 24% He
 - 1% rest ("metals")
- ⇒ expect fundamentally different internal structure!

Interiors 7



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Structure: Gas Giants

Structure of a gas giant from equation of hydrostatic equilibrium:

$$\frac{\mathrm{d}P}{\mathrm{d}r} = -\rho(r)\frac{GM(r)}{r^2}$$

To solve, need to know $\rho(r)$, $M(r) \Longrightarrow$ complicated, but doable if properties of material are known.

To guesstimate the central pressure, one can show for a planet of radius R:

$$P_{\text{central}} = \frac{2\pi}{3} G \langle \rho \rangle^2 R^2$$

Plug in numbers for Jupiter: $R = 70000 \, \mathrm{km}, \langle \rho \rangle = 1.3 \, \mathrm{g \, cm^{-3}}, \, \mathrm{get}$

 $P_{\text{central}} = 1.2 \times 10^{12} \, \text{Pa} \, (10 \times \text{Earth}).$

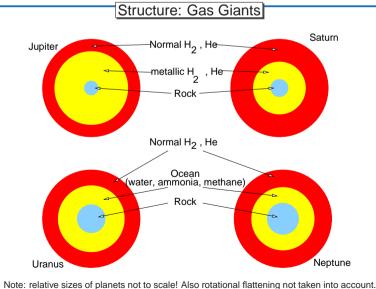
At this pressure: existence of metallic hydrogen (i.e., electrons can move freely around).

More detailed computations: metallic hydrogen from 14000-45000 km away from center

Interiors 8



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Interiors 9

here, r is now the radial distance from the planetary centre. In contrast to atmospheres, the acceleration g depends on the position, g = g(r). It is easy to show that g(r) = g(r) + g(r) +

$$g(r) = \frac{GM(r)}{r^2}$$

where M(r) is the mass of the planet contained within a radius r:

estimating the structure of atmospheres

$$M(r) = \int_{0}^{r} 4\pi \rho(r) r^{2} dr$$

(interpretation: integrate over onion shells of thickness dr and density $\rho(r)$; the mass in each of these shells is $4\pi\rho(r)dr$, summing over all onion shells gives the above

To solve the equation of the hydrostatic equilibrium one needs to know the equation of state. Unfortunately, this equation of state is generally much more complicated than for gases and often only roughly known. One can estimate, however, the order of magnitude for the pressure within a planet. In order to do so, we assume that the density is the same throughout the planet, and that it equals the planet's average density $\rho(r)=\langle \rho \rangle={
m const.}$ This is o.k. to an order of magnitude. Under this assumption,

$$M(r) = (4/3)\pi r^3 \langle \rho \rangle$$

such that the equation of hydrostatic equilibrium reads

$$\frac{\mathrm{d}P}{\mathrm{d}r} = -\langle \rho \rangle^2 G(4/3) \pi r$$

Differential equations looking like this are called separable. They can be solved "separation of variables", as we already did when computing the structure of an isothermal

First integrate both sides of the equation from r=0 to the surface of the planet at r=R

$$\int_{a}^{R} \frac{dP}{dr} dr = - \int_{a}^{R} \langle \rho \rangle^{2} G(4/3) \pi r dr$$

To integrate the left hand side of the equation, substitute $r \longrightarrow P(r)$ where P(r) is an unknown function (the pressure as a function of radius r). Luckily enough, we only need to know its values at r=0 and r=R (the "boundary conditions"). By definition of the surface of the planet, the pressure at r=R will be P(R)=0 to very good

$$\int_{0}^{R} \frac{dP}{dr} dr = P(R) - P(0) = -P(0) =: -P_{c}$$

6-50

The right hand side of the equation is easily found as well:

$$-\int_0^R \langle \rho \rangle^2 G(4/3) \pi r \, \mathrm{d} r = -\langle \rho \rangle^2 (4\pi/3) G \int_0^R r \, \mathrm{d} r = -\langle \rho \rangle^2 (4\pi/3) G R^2/2 = -\frac{2\pi}{3} \langle \rho \rangle^2 R^2$$

such that

$$P_c = \frac{2\pi}{2} \langle \rho \rangle^2 F$$

As a rule of thumb, this formula gives central pressures that are correct to better than a factor of 10 compared to the detailed theory.

