

Extragalactic X-ray and Gamma sources

Active Galactic Nuclei

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1 General Characteristics

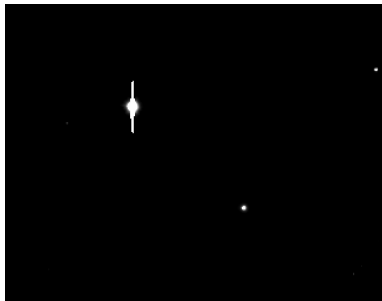
2 Classification, Unified Model

3 Energy Gain

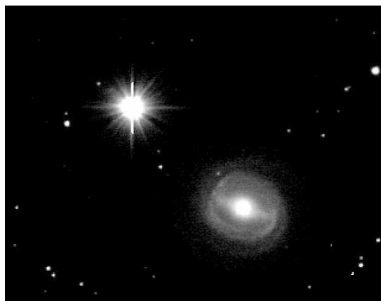
4 SED

5 Further Jet Physics

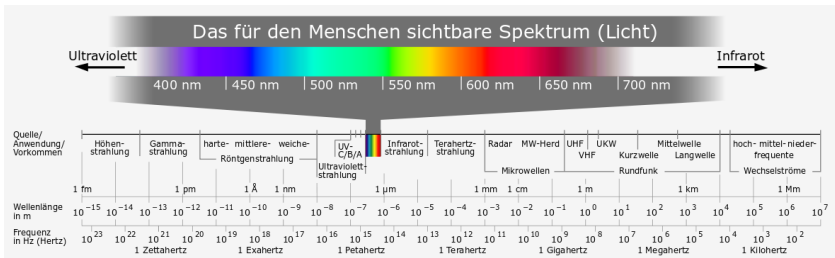
6 research



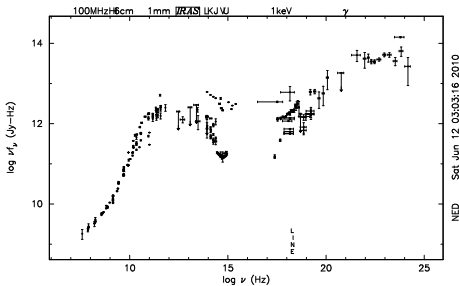
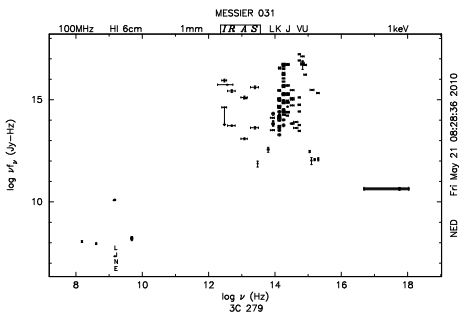
NGC3783 linear intensity scale



NGC3783 logarithmic intensity scale



Wikipedia

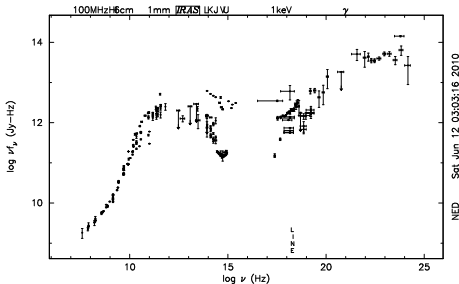
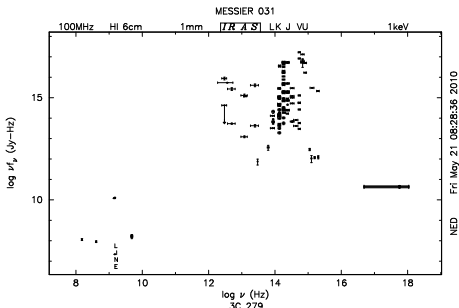


normal galaxies, e.g. M31 - Infrared:

- ▶ concentrated thermal emission at IR and optical wavebands
- ▶ diameter ~ 43 kpc
- ▶ integrated luminosity $\sim 10^{44} \frac{\text{erg}}{\text{s}}$

AGN:

- ▶ broad, mainly non-thermal continuum emission
- ▶ diameter $\sim \text{pc}$
- ▶ integrated luminosity $\sim 10^{42}-10^{48} \frac{\text{erg}}{\text{s}} \approx 10^{10} L_{\odot}$

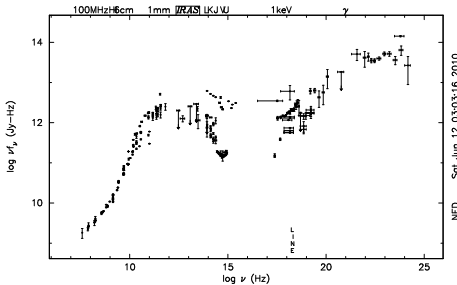
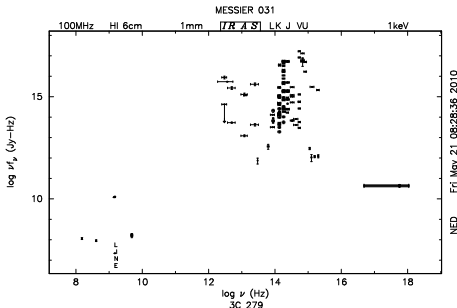


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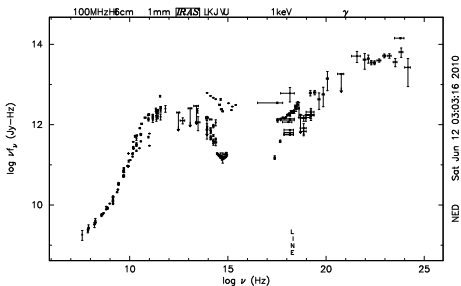
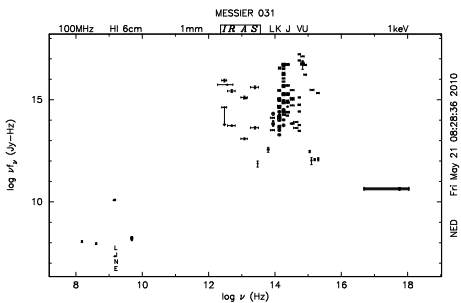


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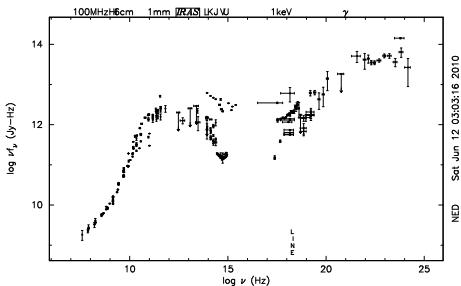
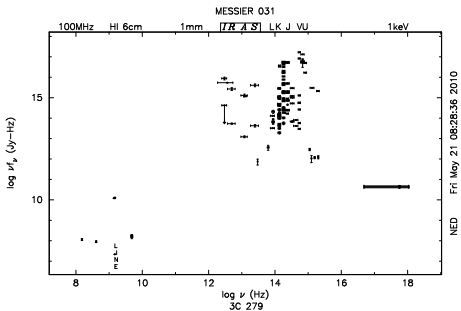


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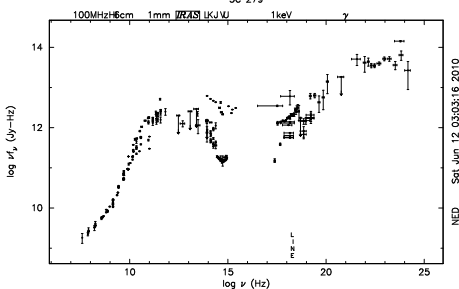
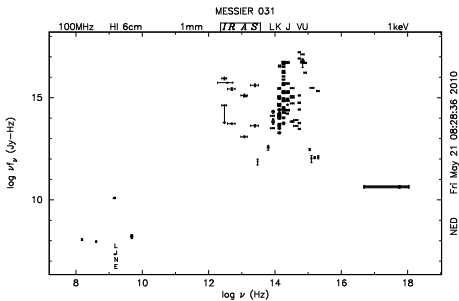


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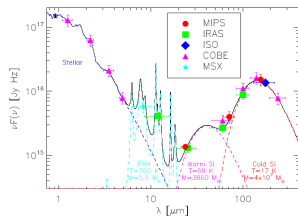
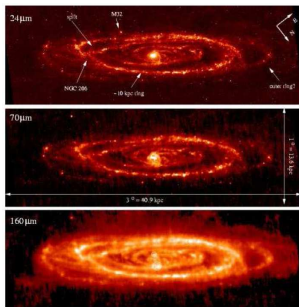
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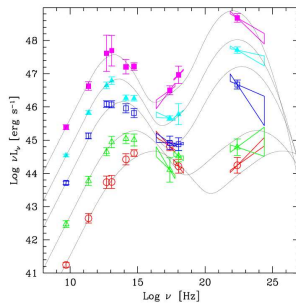
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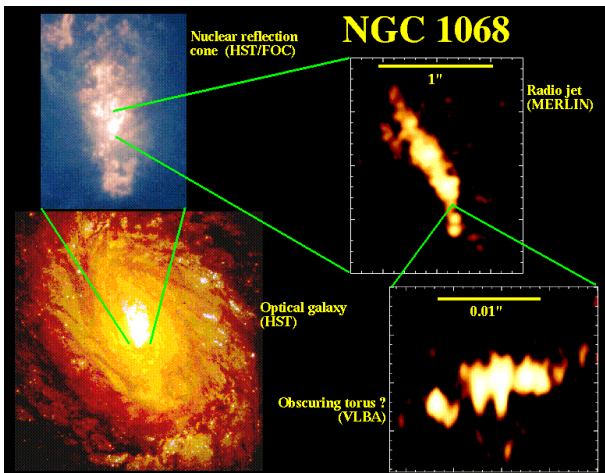
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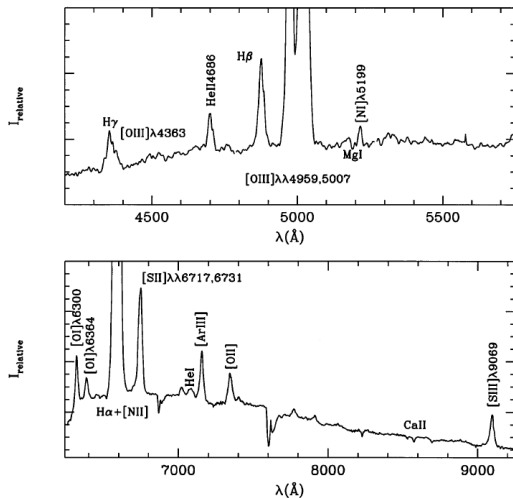
M31



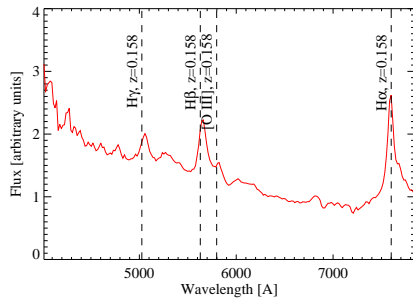
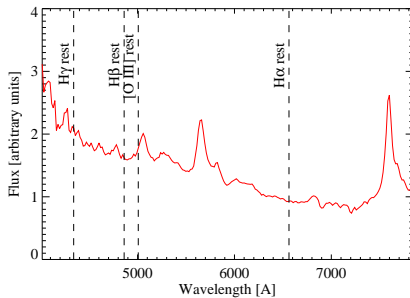
averaged SED of many blazars



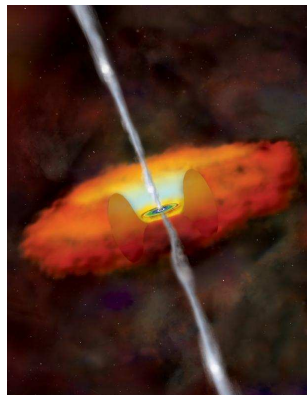
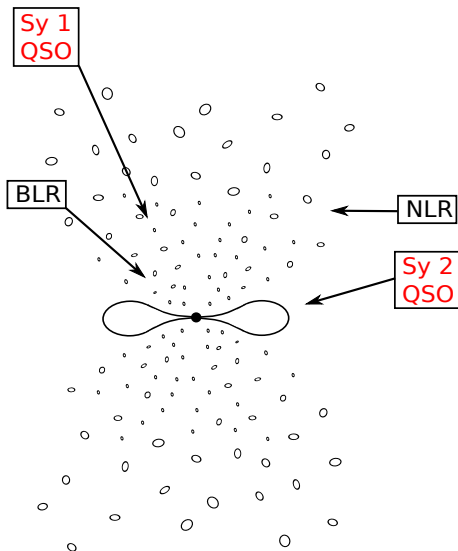




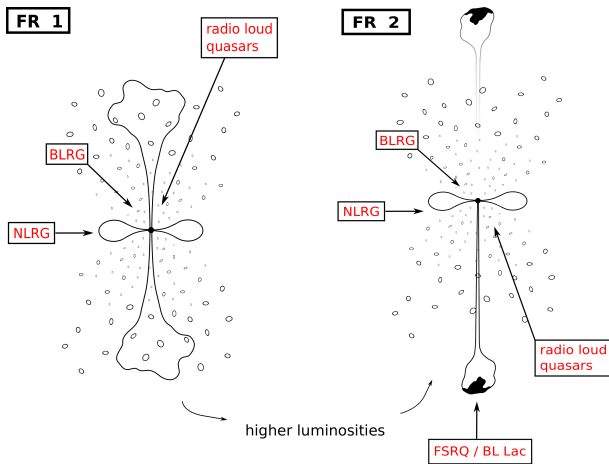
Optical spectrum of the central region of NGC 1068. Fath (1908): comparable to planetary nebula spectra, but with broad emission lines



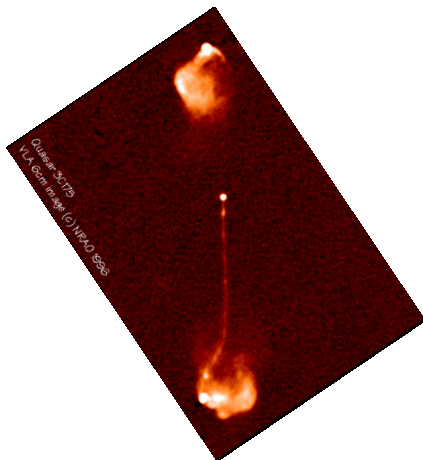
Maarten Schmidt (1962): redshift of lines \Rightarrow distance using Hubble's law $v = HD \Rightarrow$ absolute magnitude over distance modulus \Rightarrow luminosity by comparing M_{abs} with M_{\odot}
 $\Rightarrow L_{\text{quasar}} \approx 50 \cdot L_{\text{brightest galaxy}} = 4.8 \cdot 10^{12} L_{\odot}$ for 3C 273



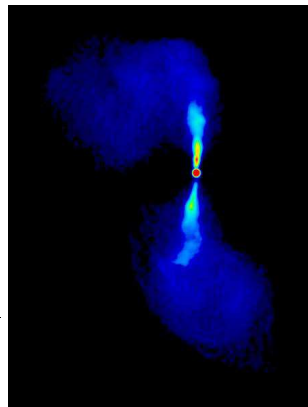
Unified model for radioquiet AGN



Unified model for radioloud AGN



FR II Type, Quasar 3C175, VLA image
at 6cm
(<http://www.cv.nrao.edu/~abridle/3c175.htm>)



FR I Type, 3C271.1, M84
(<http://nedwww.ipac.caltech.edu/>)

energy source

Which process of gaining energy is the most efficient one?

Nuclear Fusion

$$E = \epsilon mc^2 \quad (1)$$

$L \approx 10^{47}$ erg/s over 10^7 yrs ($\approx 3.2 \cdot 10^{61}$ erg) requires:

$$m = \frac{E}{\epsilon c^2} \approx 2.2 \cdot 10^9 M_{\odot} \quad (2)$$

$(1 - \epsilon)m \Rightarrow$ "fusion-waste"!

Schwarzschildradius of that mass:

$$r_s = \frac{2Gm}{c^2} \approx 6.6 \cdot 10^{12} \text{ m} \quad (3)$$

$\rightarrow \epsilon = 0.008 \Rightarrow$ energy yield $\approx 7.2 \cdot 10^{18}$ erg/g

Gravitation $\Rightarrow \epsilon \approx 0.1 \Rightarrow$ energy yield $\approx 10^{20}$ erg/g

accretion process I

optical thick accretion disc

⇒ balance between **radiation** and **gravitation**

- ▶ angular momentum → no accretion
- ▶ frictional force $F_{\text{fr}} \ll F_{\text{grav}} \Rightarrow$ Kepler orbits
- ▶ differential rotation \Rightarrow heating \Rightarrow outward loss of angular momentum \Rightarrow accretion

Radiation

$$\Delta E = \frac{GM_{\bullet}m}{r} - \frac{GM_{\bullet}m}{r + \Delta r} \approx \frac{GM_{\bullet}m}{r^2} \Delta r \quad (4)$$

virial theorem: $E_{\text{kin}} = -1/2E_{\text{pot}} = -1/2\Delta E$

$\Delta E - E_{\text{kin}} = E_i \Rightarrow$ heating \Rightarrow radiation

Using $[L] = \text{erg/s}$:

$$\Delta L = \frac{GM_{\bullet}\dot{m}}{2r^2} \Delta r \quad (5)$$

with **accretion rate** \dot{m} .

Eddington Luminosity

condition for matter being accreted (optically thick discs)

$$\frac{dp_{grav}}{dr} \stackrel{!}{>} \frac{dp_{rad}}{dr} \quad (6)$$

⇒ upper luminosity (**Eddington Luminosity** L_{Edd} , see handout)

$$L \stackrel{!}{<} L_{Edd} = \frac{GM_{\bullet} m_{HC}}{\sigma_T} \approx 1.3 \cdot 10^{38} \text{ erg/s} \cdot \frac{M_{\bullet}}{M_{\odot}} \quad (7)$$

upper **accretion rate** \dot{M}_{Edd} :

$$L_{Edd} = \eta \dot{M}_{Edd} c^2 \quad (8)$$

$$\Rightarrow \dot{M}_{Edd} = \frac{L_{Edd}}{\eta c^2} \approx 2 M_{\odot}/\text{yr} \quad (9)$$

With an efficiency η of ≥ 0.12 due to high optical depth as “resistance” for photons.

Temperature Profile

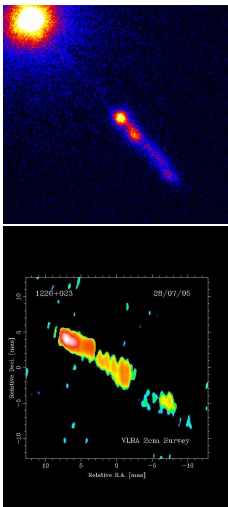
Black body radiation (optical thick) → temperature layer with **Planck Law**

$$\Delta L = 4\pi r \Delta r \sigma T^4$$

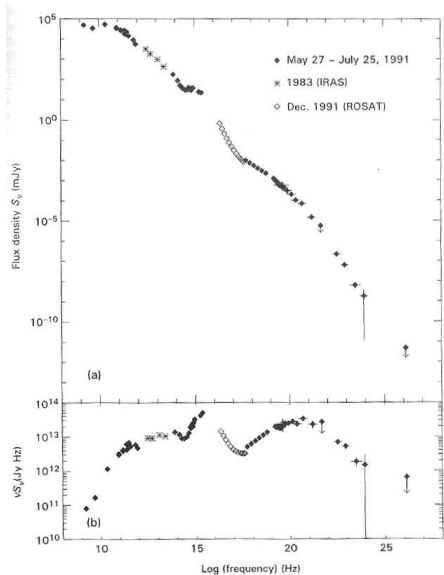
$$T(r) = \left(\frac{L}{4\pi r^3 \sigma_{\text{SB}}} \right)^{-1/4} = \left(\frac{GM_{\bullet} \dot{m}}{8\pi r^3 \sigma_{\text{SB}}} \right)^{-1/4} \underset{r_s = 2GM_{\bullet}/c^2}{=} \left(\frac{c^6}{64\pi \sigma_{\text{SB}} G^2} \right)^{1/4} \dot{m}^{1/4} M_{\bullet}^{-1/2} \left(\frac{r}{r_s} \right)^{-3/4} \quad (10)$$

→ r fixed, $\dot{m} \uparrow \Rightarrow T \uparrow$

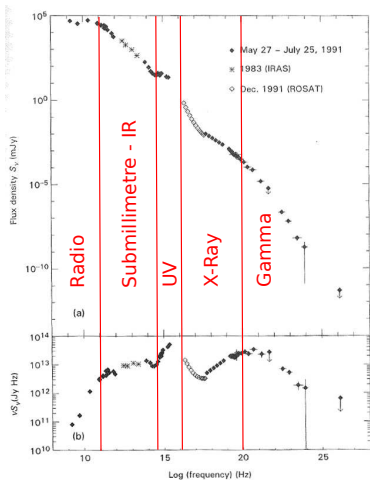
→ $M_{\bullet} \uparrow$, reached temperatures ↓



top: 3C273 in X-Rays (NASA/CXS/SAO, 2003), bottom: Jet of 3C273 in 2cm, VLBA (NRAO, Kellermann 1998)



Spectral Energy Distribution (SED) of 3C273 ([13], p.149)



- ▶ need many instruments on earth and in orbit measuring “simultaneous” if possible

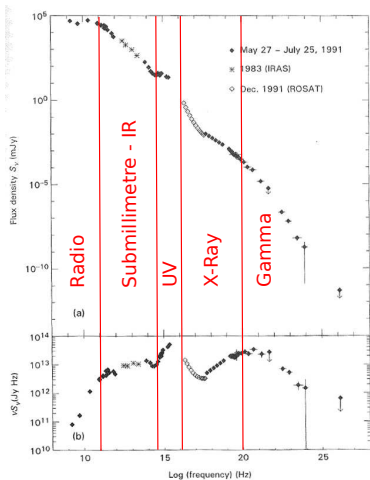
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- ▶ units: $[\nu S_\nu] = \text{W}/\text{m}^2$

$$L = \int_{\nu_1}^{\nu_2} S_\nu d\nu = \int_{\ln \nu_1}^{\ln \nu_2} \nu S_\nu d \ln \nu$$

- ▶ $(\log S_\nu - \log \nu)$: equal energy at all frequencies \rightarrow spectrum with $\alpha = -1$
- ▶ $(\log \nu S_\nu - \log \nu)$: equal energy at all frequencies \rightarrow flat spectrum with $\alpha = 0 \Rightarrow$ good indicator for above-average flux (bumps...)
- ▶ overall radiation follows powerlaws like $S_\nu \sim \nu^{-\alpha}$



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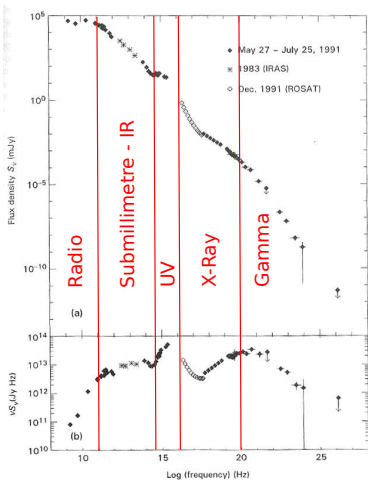
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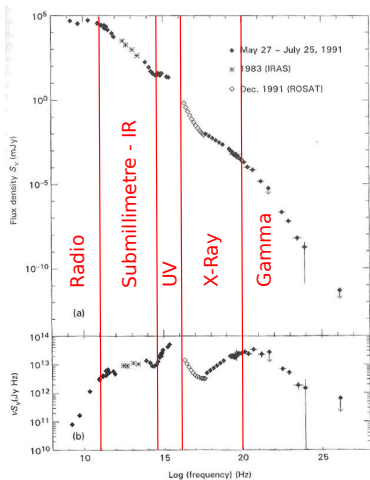
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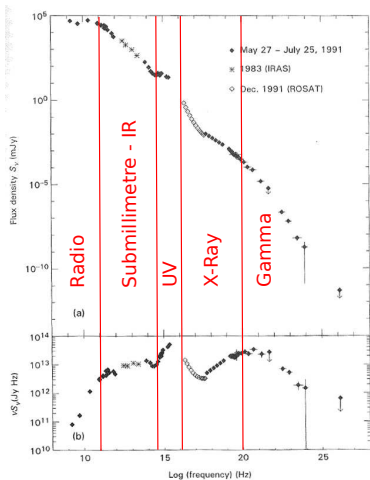
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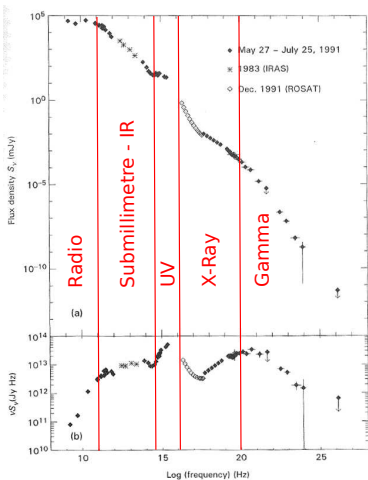
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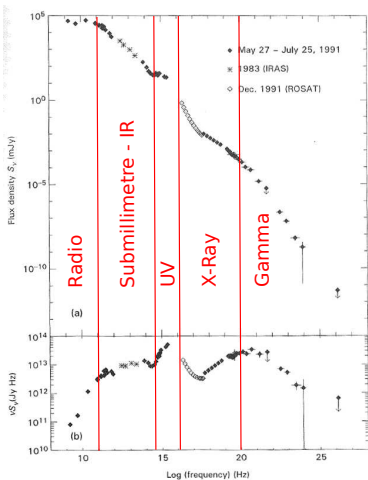
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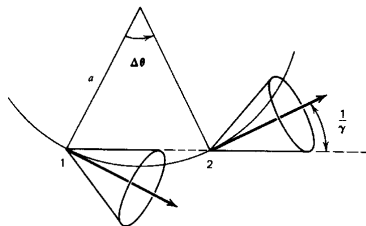


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Radio - synchrotron emission (circular orbits)



→ high degree of linear **polarization**.

→ Electron **frame of rest: radial symmetric emission**

whole emitted power:

$$I = \frac{4}{3} \sigma_T c \gamma^2 U_{\text{mag}} \quad (11)$$

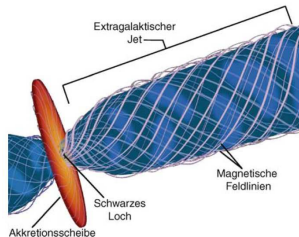
“**cooling time**” (energy decreased by factor 2):

$$t = \frac{3}{2} \frac{m^4 c^7}{e^4 B^2 E_0} \quad (12)$$

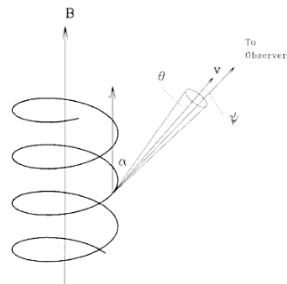
Radio - synchrotron emission (realistic conditions)

less massive particles (electrons) \Rightarrow most efficient energy loss \Rightarrow seem to form a leptonic plasma

More realistic conditions (see handout):



theoretical model of an AGN



helical trajectory of electrons around H-fieldlines

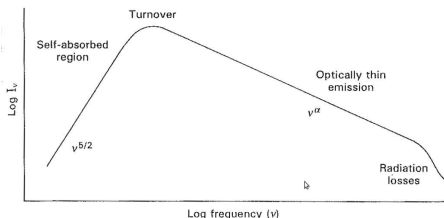
Radio - synchrotron emission (ensemble of electrons)

→ powerlaw-distribution of electrons in jet plasma: $N(E)dE \sim E^{-s}dE$

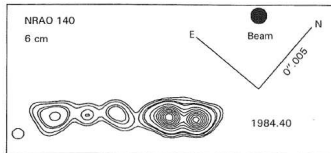
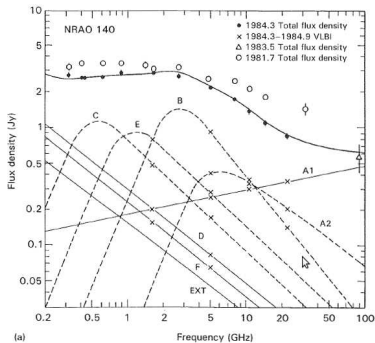
gained energy by radiation = **lost energy** through emission
 = **particle distribution** · **synchrotron emission**

$$I_\nu d\nu = \eta(E)dE = N(E)dE \cdot \frac{dE}{dt}$$

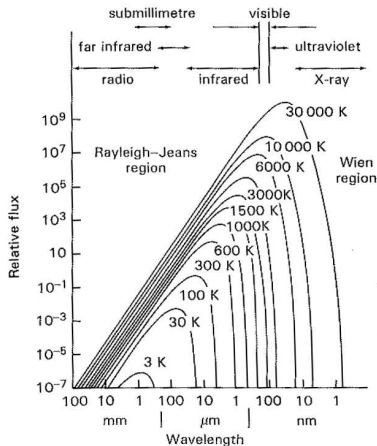
$$I_\nu \sim \begin{cases} B^{-1/2} \nu^{5/2} & \nu < \nu_c & \text{synchrotron self absorption} \\ \nu^{-(p-1)/2} & \nu > \nu_c & \text{optical thin} \end{cases} \quad (13)$$



Radio - synchrotron emission (ensemble of electrons)



Submillimetre - IR (thermal black body radiation)

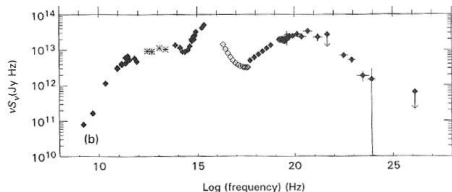


- ▶ most likely: thermal BB emission from central parts of AGN (Torus, gas, dust)
- ▶ temperatures for dust: $\approx 20 - 80$ K
- ▶ **no polarization** \rightarrow **thermal emission!**
- ▶ Planck's law:

$$B_\nu(T) = \frac{8\pi h\nu^3}{c^3 e^{\frac{h\nu}{kT}} - 1} \quad (14)$$

black body emission at different temperatures

IR - UV (thermal black body radiation)



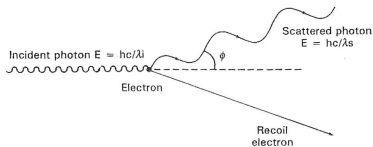
IR Bump

- ▶ near $\approx 10^{13}$ keV
- ▶ thermal emission of warm dust ($T > 2000$ K) near black hole

UV Bump

- ▶ in general: strong, broad line emission from BLR/NRL \rightarrow continuum more difficult to model than in IR!
- ▶ Big Blue Bump (**BBB**) from hot accretion disc or free-free emission (bremsstrahlung)
- ▶ thermal BB emission of the temperature-profile

X-Ray - Compton scattering



Compton scattering: energy transfer
photon → **electron**
inverse Compton scattering: energy
 transfer **electron** → **photon**

$$\lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta) \quad (15)$$

$$E'_e = \frac{E}{1 + \frac{E}{m_e c^2} (1 - \cos \theta)} \quad (16)$$

$$\frac{\Delta E}{E} \approx -\frac{E}{m_e c^2} \quad (E \ll m_e c^2) \quad (17)$$

From Eq.16 for many scattering events (cf. Eq.11):

$$I = \frac{dE}{dt} = \frac{4}{3} \sigma_T c \gamma^2 U_{\text{el}} \quad (18)$$

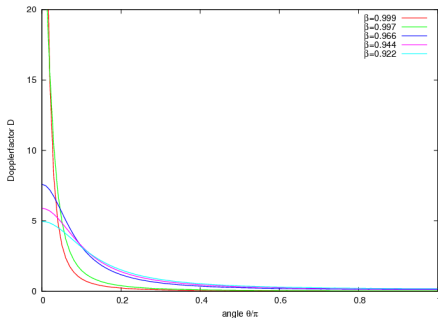
relativistic boosting

Time dilatation causes **relativistic Doppler effect** with

$$v_{\text{obs}} = \frac{v_{\text{em}}}{\gamma(1 - \beta \cos \theta)} \quad (19)$$

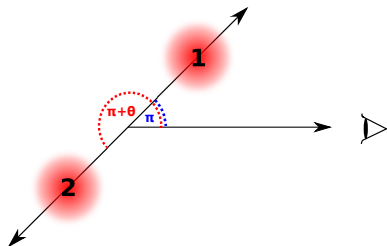
with the relativistic **Doppler factor**

$$\mathcal{D} = \frac{\sqrt{1 - \beta^2}}{(1 - \beta \cos \theta)} \quad (20)$$



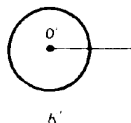
relativistic boosting - jet one sidedness

One can show, that $\frac{l_v^{\text{obs}}}{v_{\text{obs}}^3} = \frac{l_v^{\text{em}}}{v_{\text{em}}^3} \Rightarrow l_v^{\text{obs}} = \mathcal{D}^3 l_v^{\text{em}}$
 power law $l_v \sim A v^\alpha \Rightarrow l_v^{\text{obs}} = \mathcal{D}^{3-\alpha} l_v^{\text{em}}$

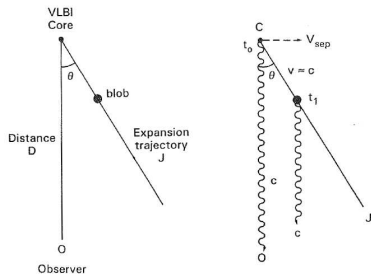


$$\frac{l_1}{l_2} = \left(\frac{1 + \beta \cos \theta}{1 - \beta \cos \theta} \right)^{3-\alpha} \quad (21)$$

In addition: **relativistic aberration**



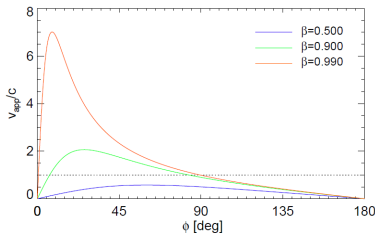
Superluminal Motion



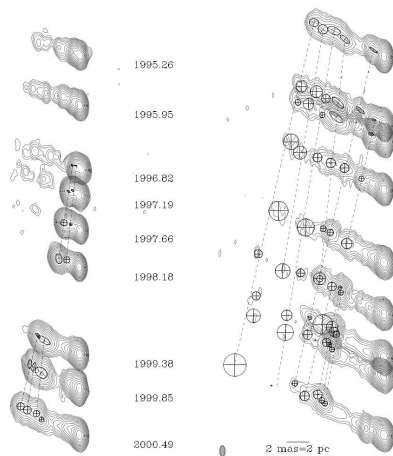
apparent speed of a blob:

$$v_{\text{app}} = \frac{v \sin \theta}{1 - \beta \cos \theta} \quad (22)$$

“superluminal” only for relativistic blob-speeds (large β) at small viewing angles Φ

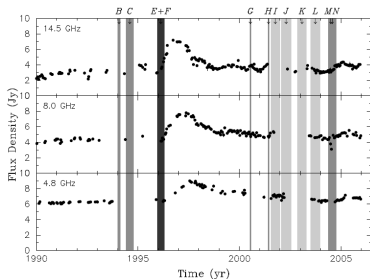


tracking flares of 3C111



VLBA monitoring at 2cm

tracking flares of 3C111

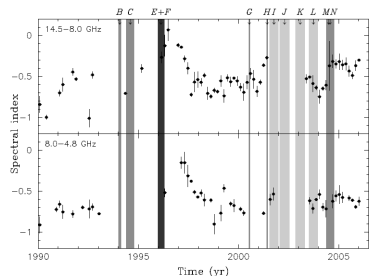


radio lightcurve (top figure)

- ▶ blob first visible at high frequencies (synchrotron self absorption mainly at lower frequencies)
- ▶ blob expands \Rightarrow less dense electron plasma \Rightarrow less synchrotron self absorption

spectral indices (bottom figure)

- ▶ spectral indices α from $I_\nu \sim \nu^\alpha$
- ▶ compact blobs in plateau-state \Rightarrow flat radio spectrum ($\alpha \approx 0$)
- ▶ decay state: blob expands \Rightarrow radio spectrum steepened ($\alpha < 0$)



Relationship between frequency bands

relationship radio - gamma ([17])

- ▶ comparison **radio** (22 GHz and 37 GHz, Metsähovi Obs.) - **gamma** (EGRET)
- ▶ **radio** emission several month **after gamma** emission
- ▶ coupling gamma - radio: both originate in same flare \leftrightarrow gamma rays from SSC-upscattering of synchrotron seed photons (from accelerated relativistic electrons)

relationship optical - radio ([16]: Generalized Shock Model)

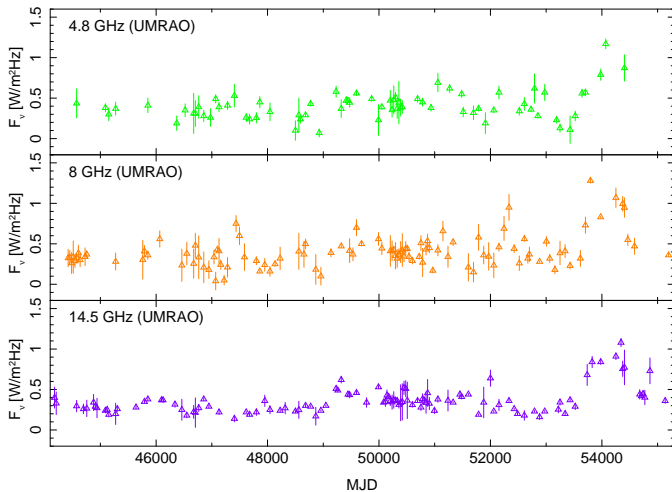
- ▶ connection: **accreted matter** (\rightarrow optical thermal emission) - **radio-flare**
- ▶ strong delay between accretion and radio-flare expected
- ▶ or: optically thin slope of synchrotron spectrum reaches optical waveband \rightarrow no delay!

UMRAO Radio Observatory

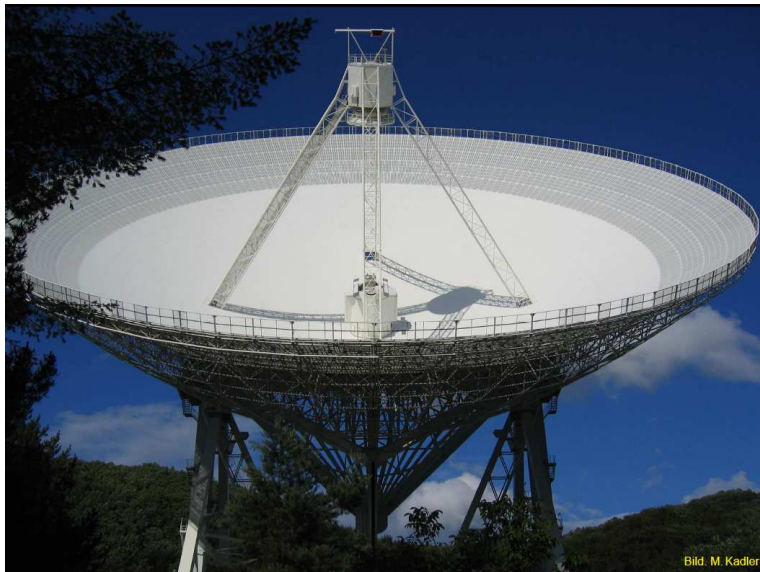


UMRAO lightcurve of PKS 2155-304

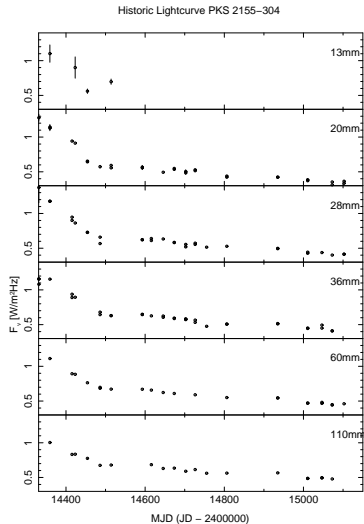
Historic Lightcurve PKS 2155-304



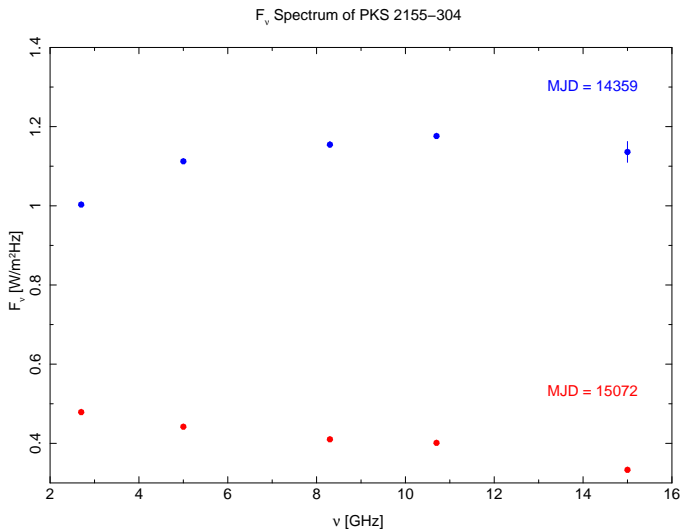
Effelsberg Radio Telescope



Effelsberg lightcurves of PKS 2155-304

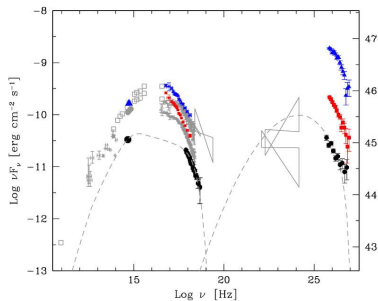


Effelsberg lightcurves of PKS 2155-304



Multiwavelength observations of the 2006 flare of PKS 2155-304

simultaneous observation of the flare with HESS (gamma), Chandra (X-ray), RossiXTE (X-ray), Bronberg Obs. (optical)



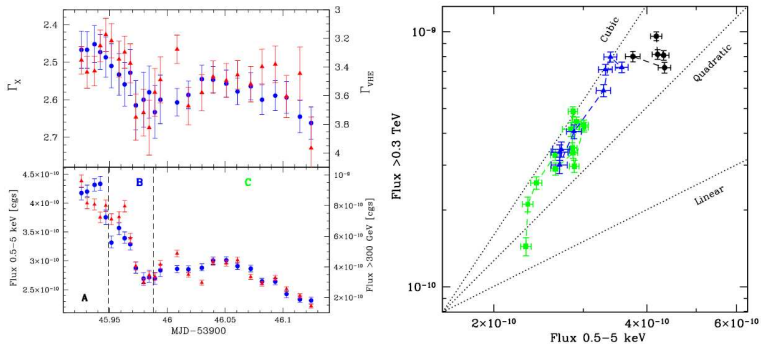
SED of PKS 2155-304 with highest and lowest states during this observation

- ▶ first peak, right slope: X-Ray (Synchrotron emission)
- ▶ second peak: Gamma (inverse Compton emission)
- ▶ flare not moving though frequencies with time

Why X-Ray through Synchrotron emission??

- ▶ **“blue blazars”**
→ less external photons → less Compton cooling of electrons → overall higher photon energies due to synchrotron or inverse Compton recoil
- ▶ **“red blazars”**
→ higher photon density → lower photon energies

Multiwavelength observations of the 2006 flare of PKS 2155-304



plot of spectral index' and flux variability

- ▶ strong correlation between X-ray and γ -ray flux (synchrotron and inverse Compton emission \rightarrow as already shown)
- ▶ γ -ray flux decreases approximately with cube of X-ray flux ($F_\gamma \sim F_X^3$)

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