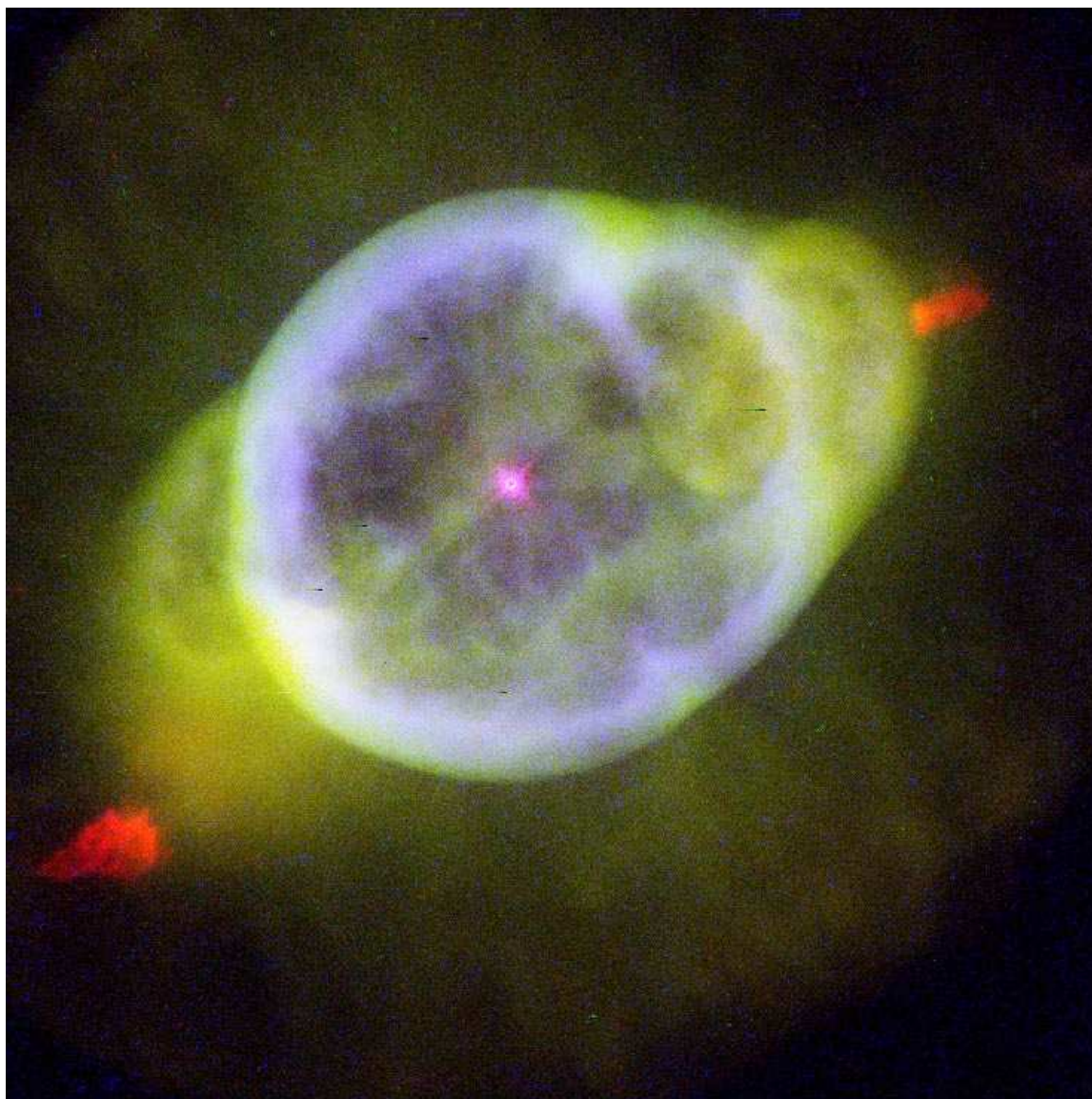


Line Diagnostics

Planetary Nebulae, I



NGC 3242 with HST

IAAT



NGC 6543

PR95-01a • ST ScI OPO • January 1995 • P. Harrington (U.MD), NASA

HST • WFPC

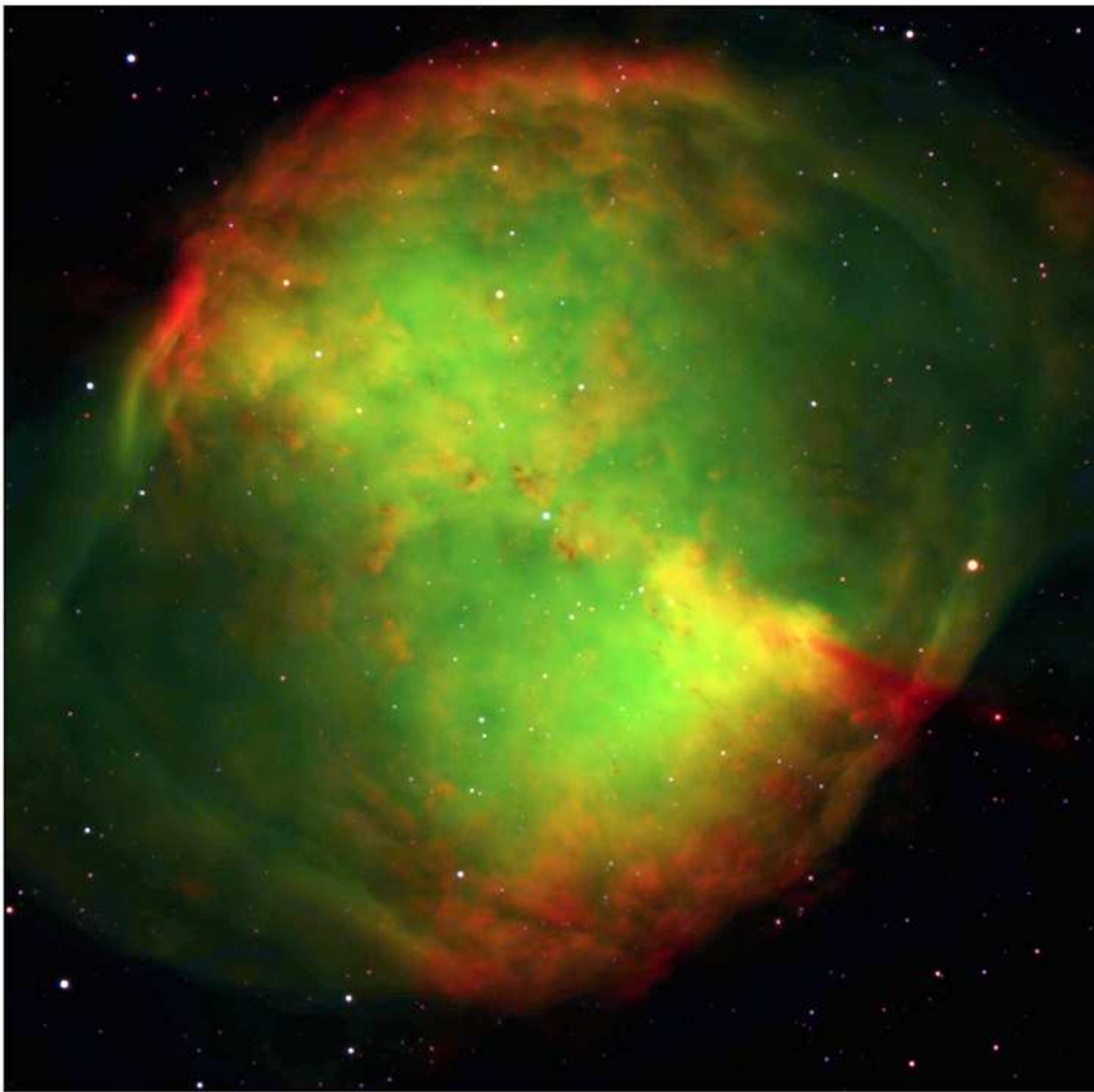
12/13/94

Planetary Nebula IC 418



Hubble
Heritage

PRC00-28 • NASA and The Hubble Heritage Team (STScI/AURA) • HST/WFPC2



Planetary Nebula NGC 6853 (M 27) - VLT UT1+FORS1

ESO PR Photo 38a/98 (7 October 1998)

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Blue: broad band blue filter, green: O II, red: H α

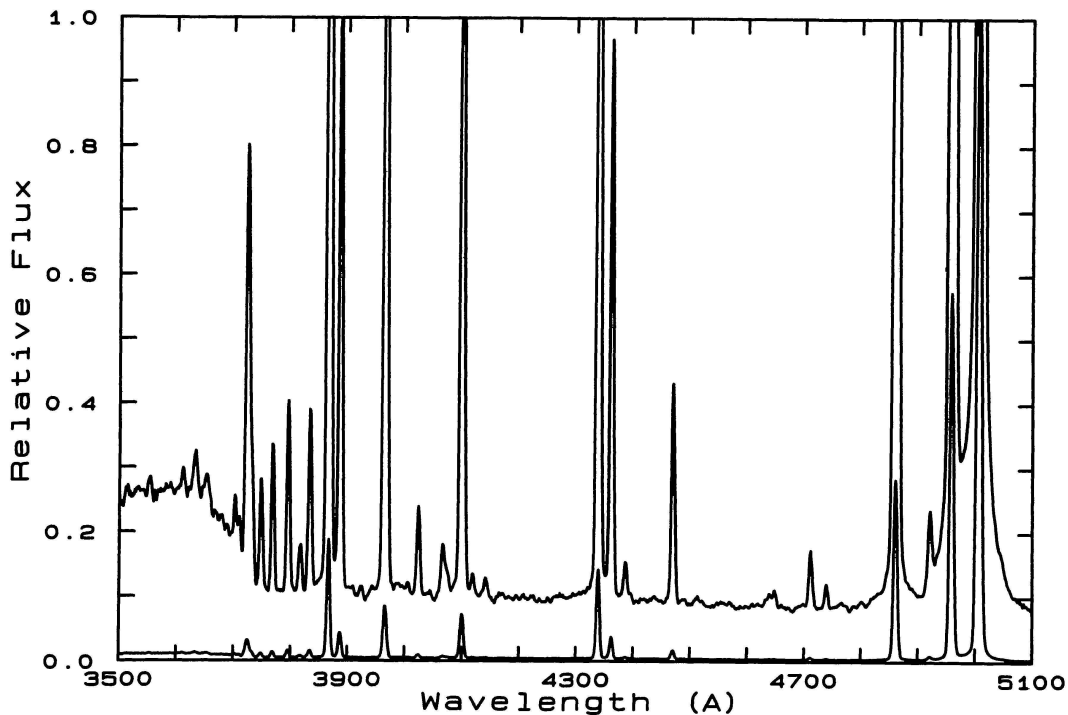


FIG. 1a

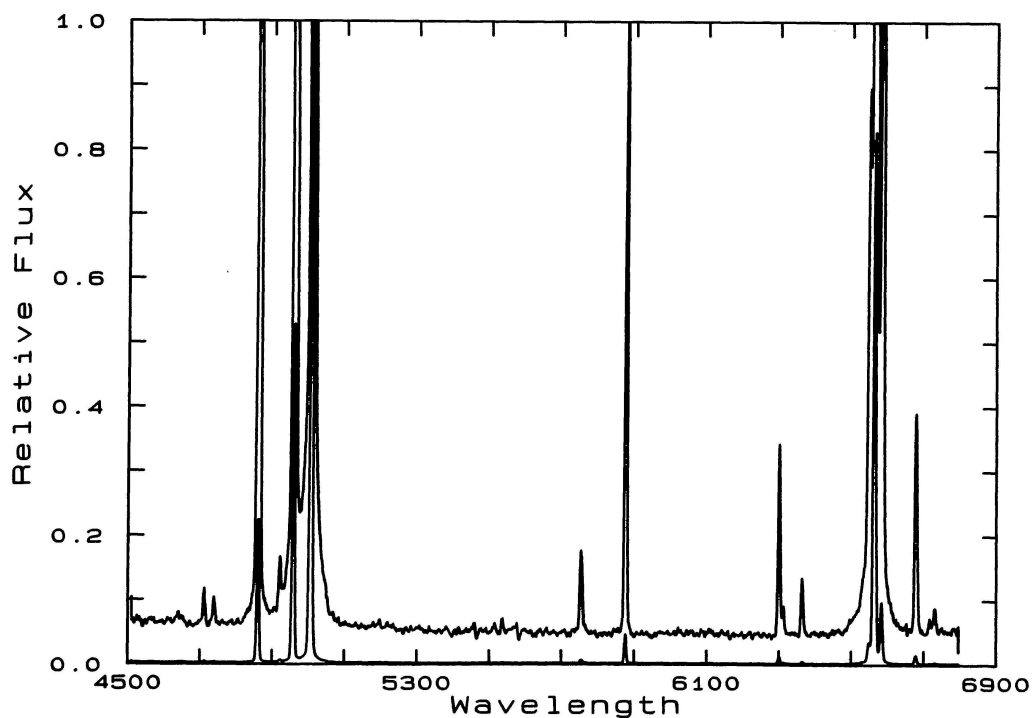
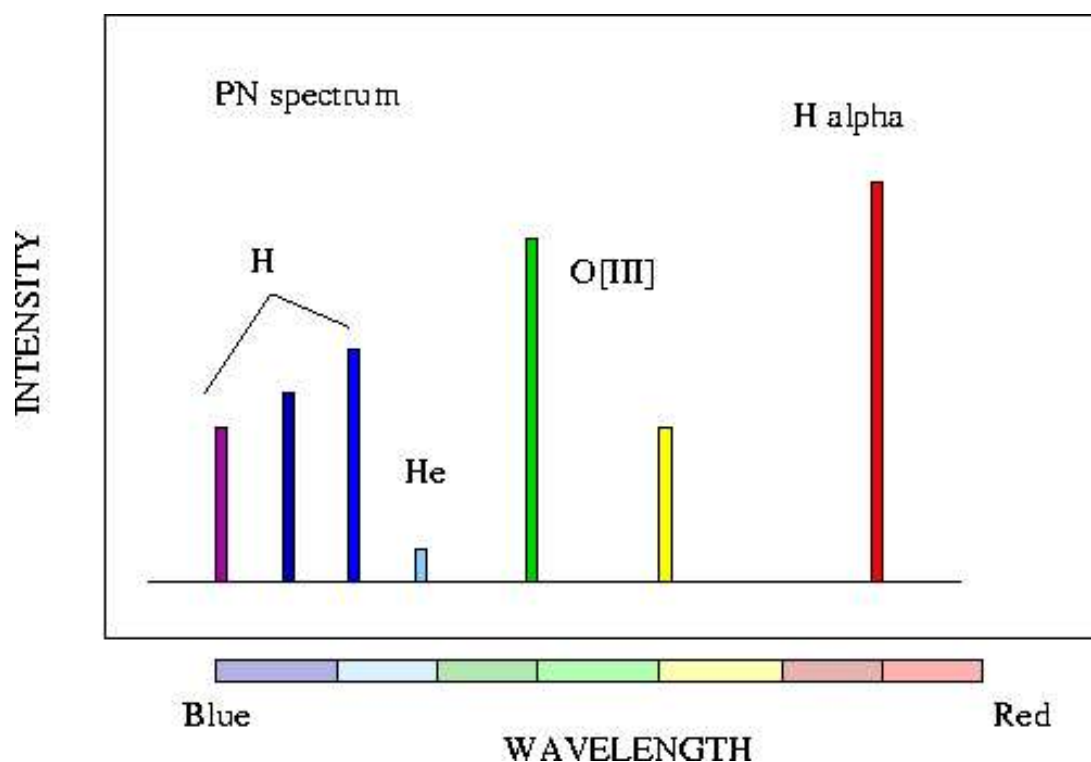


FIG. 1b

FIG. 1.—Spectral scans of the low-excitation planetary NGC 6833. All scans are corrected for effects of atmospheric extinction, response functions, and nonlinearity of dispersion. Nebular fluxes were measured in units of $\text{ergs cm}^{-2} \text{s}^{-1} \text{\AA}^{-1}$. We give here relative fluxes. All data were secured with the 3 m Shane telescope and image-tube scanner on 1981 August 7. All scans of NGC 6833 have a resolution of 9 Å. (a) The regions 3500–5100 Å. The two magnifications are 1× (scaled to strongest line $\lambda 5007$) and 25×. Note the strong background continuum. (b) The region 4500–6800 Å, magnifications 1× and 25×. (c) The region 6200 to 7800, magnifications 1× and 25×. Neither [N II] nor [S II] is prominent in this object, although the [S III] auroral-type transition $\lambda 6312$ is fairly strong.

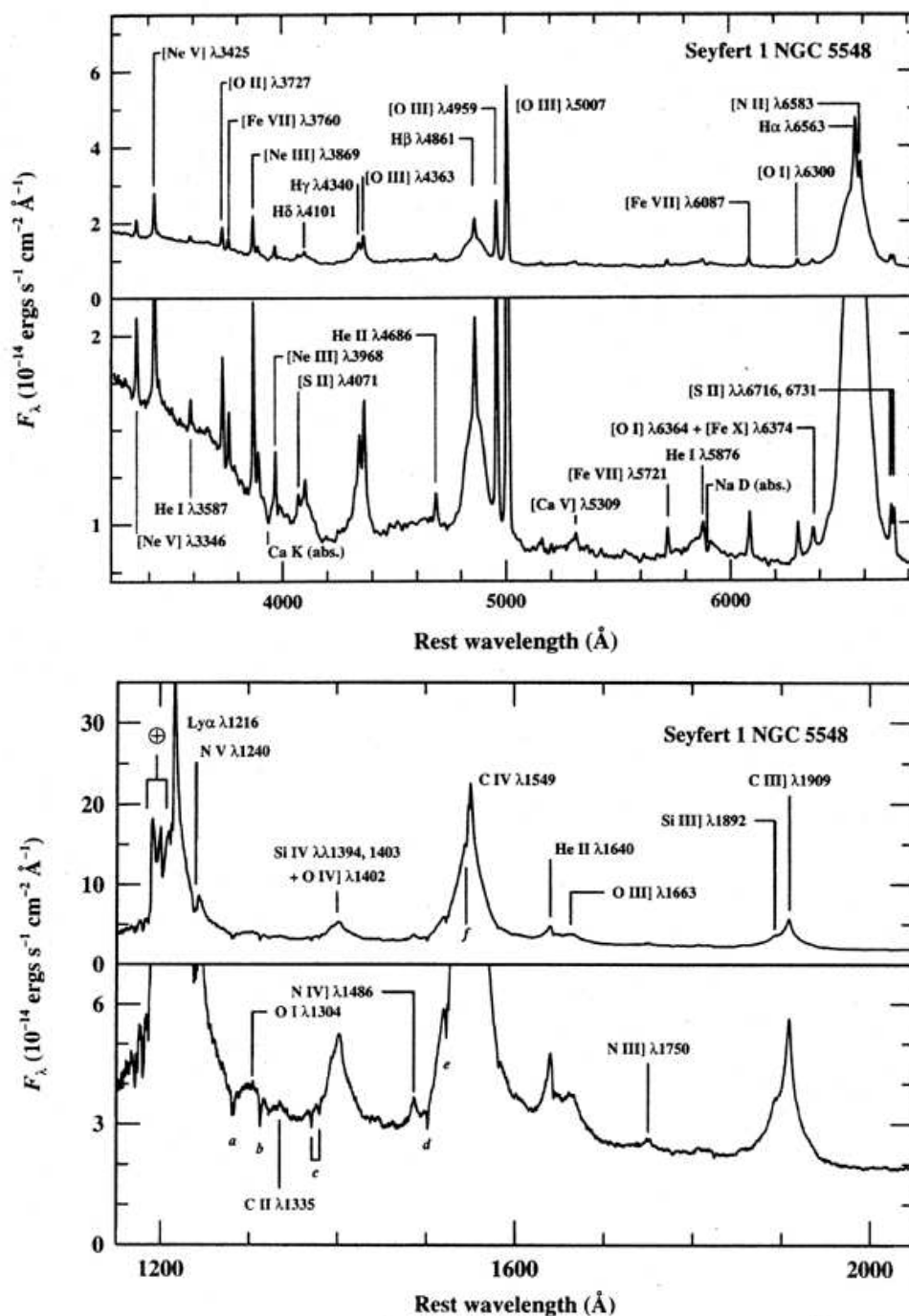
Planetary Nebulae, VI



www.ucolick.org/~bolte/AY4/notes10/node3.html

Typical spectra of planetary nebulae are dominated by Hydrogen lines, plus emission from O III at 5007 Å (“**nebulium**”).

Seyfert Galaxies



Peterson (1997)

Optical/UV spectrum of NGC 5548: a typical **Seyfert 1 Galaxy**.

Photoionization, I

Assume: cloud irradiated by photons

Goal: only source for ionization: **photoionization**

Equilibrium: number ionizations = number of recombinations

$$\int_{\nu_{\text{ion}}}^{\infty} a(\nu) \frac{F_{\nu}}{h\nu} n(X^r) d\nu = \alpha(T) n_e n(X^{r+1}) \quad (7.1)$$

where

$a(\nu)$: photoionization cross section (cm^2 ; $\propto E^{-3}$)

$\alpha(T_e)$: Recombination coefficient ($\text{cm}^3 \text{s}^{-1}$)

n_i : particle density (cm^{-3})

F_{ν} : local photon flux, $\text{erg cm}^{-2} \text{s}^{-1} \text{keV}^{-1}$,

$$F_{\nu} = \frac{L_{\nu}}{4\pi D^2} \quad (7.2)$$

Since $a(\nu)$ quickly decreasing function:

$$\frac{n(X^r)}{n(X^{r+1})} \sim \frac{a(\nu_{\text{ion}})}{\alpha(T)} \frac{L}{4\pi D^2 n_e} \frac{1}{h\nu_{\text{ion}}} \quad (7.3)$$

i.e., ionization equilibrium mainly depends on

$$U = \frac{L/4\pi D^2 h\nu_{\text{ion}}}{n_e} \frac{1}{c} = \frac{\# \text{ ionizing photons/cm}^3}{\# \text{ electrons/cm}^3} \quad (7.4)$$

the **ionization parameter**

many other definitions available!

Photoionization, II

In reality, other physical processes need to be considered:

Ionization:

- Photoionization
- collisional Ionization
- Auger-Ionization

Recombination:

- radiative recombination
- dielectric recombination

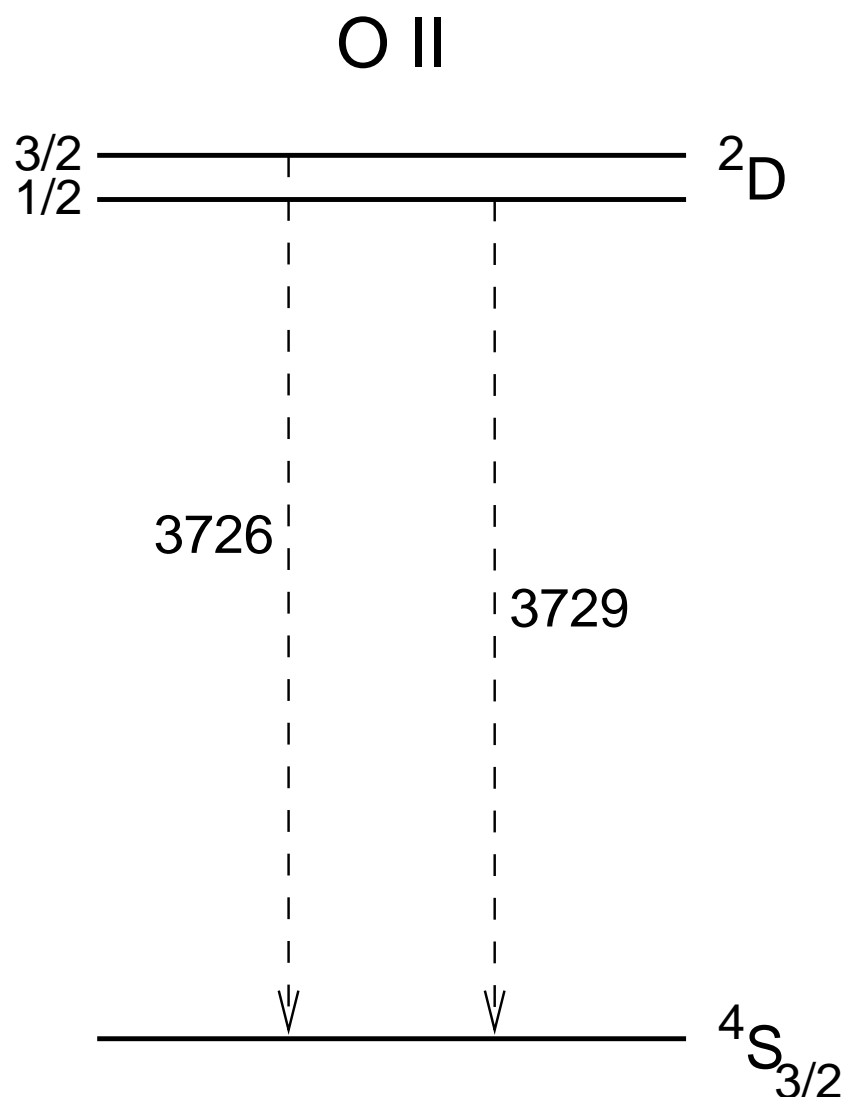
Continuum Processes:

- Bremsstrahlung
- Compton-Scattering

Real life: Solution of RT problem using advanced radiation codes such as [Cloudy](#) or [XSTAR](#) (*not* worthwhile to develop your own code...).

Line Diagnostics: Density, I

Choose atom with two levels with almost **same excitation energy**. Either **radiative** or **collisional deexcitation**.



For $N_e \approx 1000 \text{ cm}^{-3}$ use [O II] 3729/3726, for higher densities: C II

Line Diagnostics: Density, II

Rate equations in equilibrium:

$$n_1 n_e C_{12} = n_2 A_{21} + n_2 n_e C_{21} \quad (7.5)$$

$$n_1 n_e C_{13} = n_3 A_{31} + n_3 n_e C_{31} \quad (7.6)$$

where A_{ij} **Einstein-Coefficient**, C_{ij} coefficient for collisional (de)excitation.

Computation of C_{ij} :

For de-excitation:

$$C_{21} = \int_0^{\infty} \sigma_{21}(v) v f(v) d^3v \quad (7.7)$$

where σ : cross section, $f(v)$ Maxwell.

One can show that

$$\sigma_{21}(v) = \frac{\pi \hbar^2}{m^2 v^2} \frac{\Omega_{21}}{g_2} \quad (7.8)$$

where Ω_{21} : **collision strength**.

Therefore

$$C_{21} = \frac{\hbar^2}{m^{3/2}} \frac{\Omega_{21}}{g_2} \left(\frac{2\pi}{kT} \right)^{1/2} \sim \frac{8.616 \times 10^{-6} \Omega_{21}}{T^{1/2}} \frac{1}{g_2} \text{cm}^3 \text{s}^{-1} \quad (7.9)$$

Because of Microreversibility

$$C_{12} = \frac{g_2}{g_1} C_{21} \exp(-E_{12}/kT) \quad (7.10)$$

Line Diagnostics: Density, III

Solve rate equations

$$\frac{n_2}{n_1} = \frac{n_e C_{12}}{A_{21} + n_e C_{21}} \quad (7.11)$$

$$= \frac{n_e}{A_{21} + n_e C_{21}} \frac{g_2}{g_1} C_{21} \exp(-E_{12}/kT) \quad (7.12)$$

and a similar equation for n_3/n_1 .

Intensity of the line (assuming cloud is optically thin)

$$I_{21} = \frac{A_{21} n_2 h \nu_{21}}{4\pi} \quad (7.13)$$

Therefore

$$\frac{I_{21}}{I_{31}} = \frac{A_{21} n_2 h \nu_{21} / 4\pi}{A_{31} n_3 h \nu_{31} / 4\pi} \quad (7.14)$$

since $\nu_{21} \sim \nu_{31} \dots$

$$= \frac{A_{21} n_2}{A_{31} n_3} \quad (7.15)$$

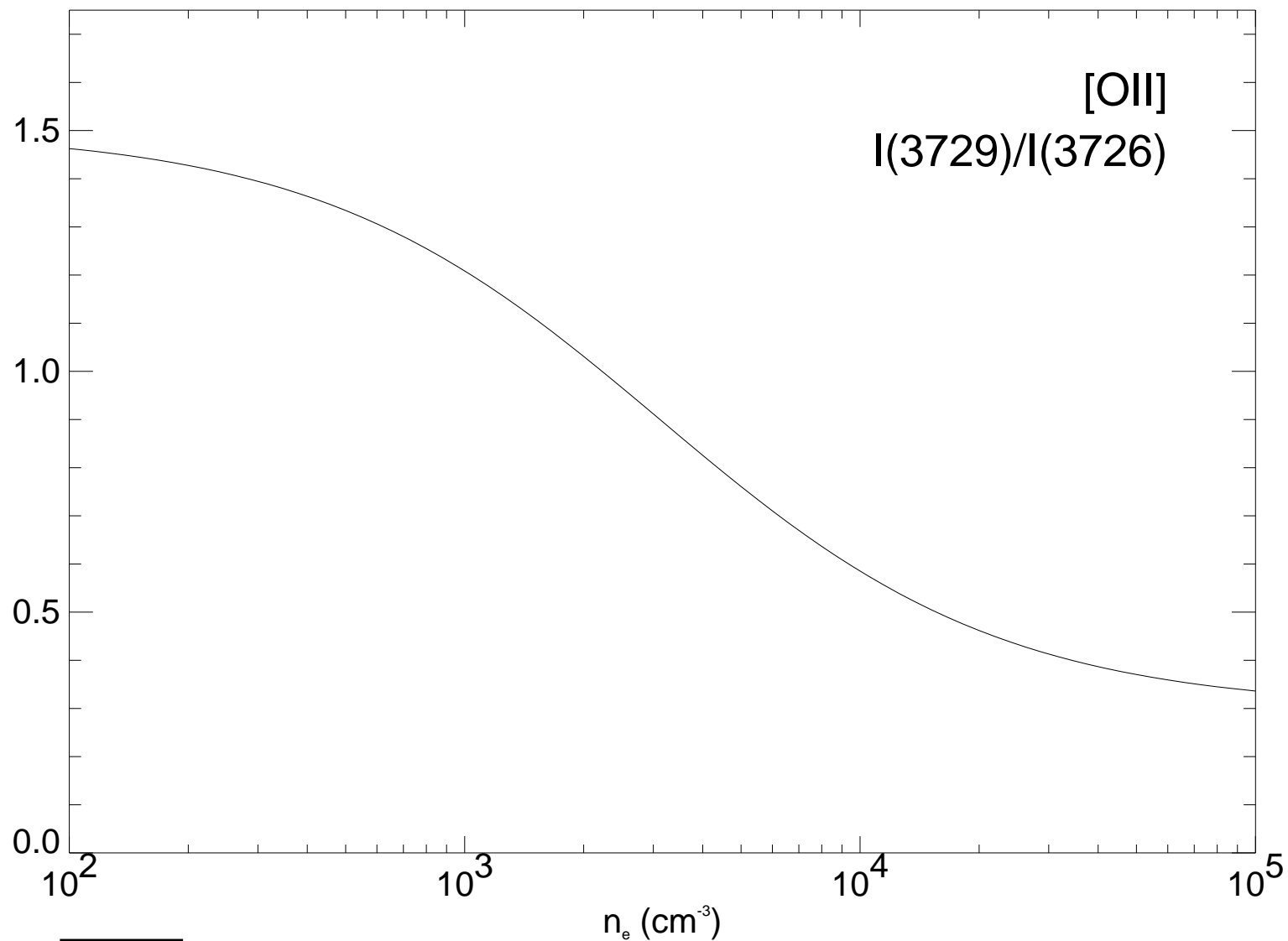
$$= \frac{C_{21} g_2 A_{21} A_{31} + n_e C_{31}}{C_{31} g_3 A_{31} A_{21} + n_e C_{21}} \exp(-E_{32}/kT) \quad (7.16)$$

$$= \frac{g_2 C_{21} (1 + n_e/n_{cr,3})}{g_3 C_{31} (1 + n_e/n_{cr,2})} \exp(-E_{32}/kT) \quad (7.17)$$

where the **critical density** is defined by

$$n_{cr,2} = \frac{A_{21}}{C_{21}} \quad (7.18)$$

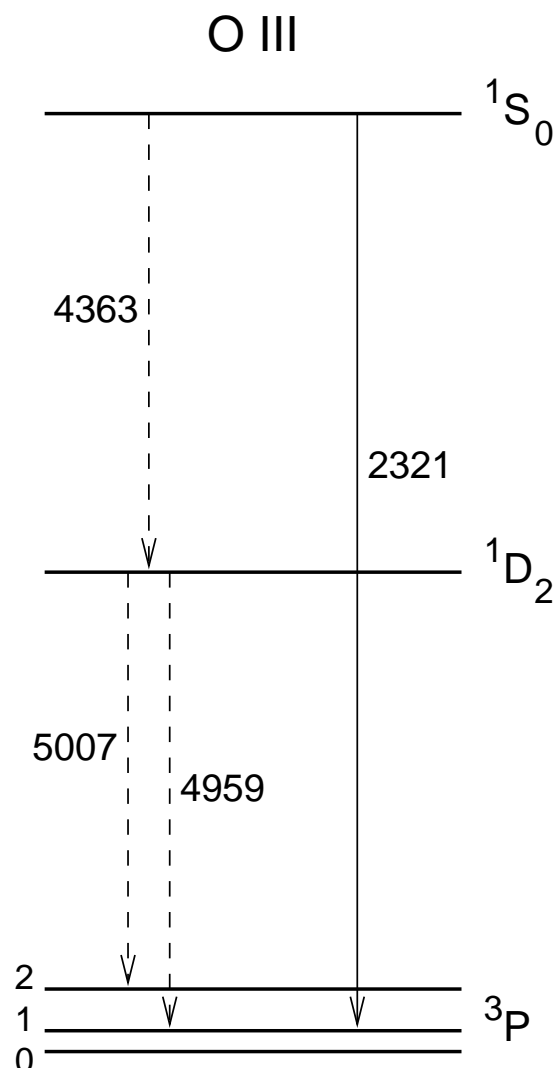
Line Diagnostics: Density, IV

**IAAT**

Line Diagnostics: Temperature, I

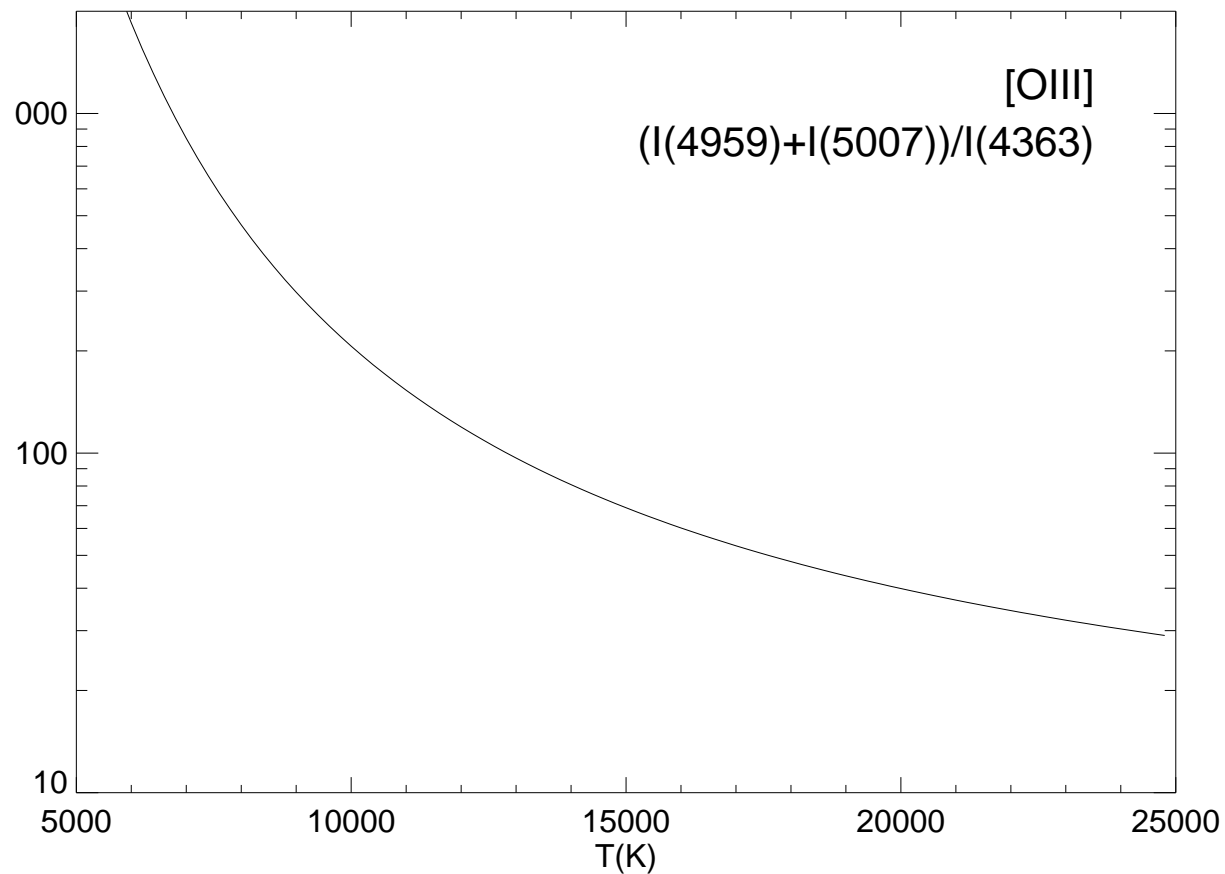
To obtain **temperature** use similar ideas. This time, use two levels with different excitation energy

⇒ Use different excitation probability of collisional excitation



For $T \sim 10000$ K, mainly O III and N III

Line Diagnostics: Temperature, II



$$\frac{I(4959 + 5007)}{4363} = \frac{7.7 \exp(3.29 \times 10^4 / T)}{1 + 4.5 \times 10^{-4} n_e T^{-1/2}}$$

IAAT

Line Diagnostics: Mass

Mass determination: Determine number of emitting atoms from line strength.

Hydrogen: $H\beta$ (less influenced by radiative transfer effects)

$$j_{H\beta} = n_e n_p \alpha_{H\beta} \frac{h\nu_{H\beta}}{4\pi} \quad (7.19)$$

$$= \frac{n_e^2}{4\pi} \alpha_{H\beta}^{\text{eff}} \frac{h\nu_{H\beta}}{4\pi} \quad (7.20)$$

$$= 1.24 \times 10^{-25} \text{ erg s}^{-1} \text{ cm}^{-3} \text{ sr}^{-1} \frac{n_e^2}{4\pi} \quad (7.21)$$

where $\alpha_{H\beta}^{\text{eff}}$: effective recombination coefficient for $n = 4 \rightarrow n = 2$ transition (weakly temperature dependent).

Total emissivity

$$L_{H\beta} = \int \int j_{H\beta} d\Omega dV \quad (7.22)$$

$$= \frac{4\pi n_e^2}{3} \cdot 1.24 \times 10^{-25} r^3 f \text{ erg s}^{-1} \propto \int n_e^2 dV \quad (7.23)$$

where $\int n_e^2 dV$: **emission measure**, and f : **filling factor**.